

## New insight into water-obsidian interaction

Y. Kudriavtsev<sup>a</sup>, R. Asomoza-Palacio<sup>a</sup>, L. Manzanilla-Naim<sup>b</sup>

<sup>a</sup>Departamento Ingeniería Eléctrica -SEES, CINVESTAV-IPN, México, DF, 07360, México.

<sup>b</sup>Instituto de Investigaciones Antropológicas, Universidad Nacional Autónoma de México.

e-mail: yuriyk@cinvestav.mx

Received 29 August 2016; accepted 23 September 2016

Hydration of obsidian samples in isotopic water  $\text{H}_2^{18}\text{O}$  was performed at temperatures from  $75^\circ\text{C}$  to  $200^\circ\text{C}$ . The depth distributions of  $^1\text{H}$  and  $^{18}\text{O}$  in the obsidians measured by SIMS were found to be drastically different. This led us to the conclusion that water molecules split into hydrogen and oxygen at the obsidian surface and then the atomic hydrogen and  $^{18}\text{O}$  isotope diffused into the obsidian via two different mechanisms. The hydrogenated obsidians were heated in vacuum without hydrogen emission. A hypothesis is put forward that the obsidian hydrogenation results from a chemical reaction between atomic hydrogen and a glass network. A new linear-parabolic equation is suggested for obsidian hydrogenation dating. Aluminosilicate glass is considered as a promising material for hydrogen fusion cells.

**Keywords:** Obsidian hydration dating; diffusion; solid state chemical reaction; SIMS; hydrogenation; central Mexico.

PACS: 61.43.Fs; 66.30.-h; 66.30.jj; 66.55.Nq; 82.80.Ms

### 1. Introduction

In 1960 Friedman and Smith [1] suggested a new method for dating of obsidian artifacts. The original obsidian hydration dating method relies upon measuring the water diffusion depth in a freshly created obsidian surface and converting thickness measurements to an age. The hydration rim can be examined in an optical microscope or by the Secondary Ion Mass Spectrometry (SIMS) method [2,3]. As an “archaeological equation” relating the artifact age  $t$  and the hydrated layer  $L$ , Friedman and Smith suggested the parabolic equation [1]:

$$L = \sqrt{(D \cdot t)}, \quad (1)$$

where  $D$  is the diffusion coefficient that depends on the temperature in the local site and on the obsidian composition.

Obsidian artifacts are widely distributed in Mesoamerica, USA, Japan, Oceania, and for this reason the obsidian hydration dating method (OHD) was met with great enthusiasm. However, it was difficult or even impossible to explain serious discrepancies between experimental facts accumulated during the last 55 years by the original OHD model. Serious discrepancies between the OHD results and the chronological data obtained by other methods were reported [4-6]. The diffusion coefficients obtained in laboratories by the intrinsic rate method [7,8] were found to differ from the coefficients calculated for archaeological artifacts by the radiocarbon calibration [9]. Some authors noted discrepancies between their experimental data and Eq. (1) [9,10]. All this gave rise to doubts in the physical basis of the OHD method [4,9,10]. Among other alternatives, penetration of atomic hydrogen rather than water into obsidian was suggested [10,11].

We performed hydration of obsidians at different temperatures and during different times in vapors of isotopic water  $^1\text{H}_2$   $^{18}\text{O}$  with 97% of the  $^{18}\text{O}$  isotope. The idea of this experiment is evident: if there is water diffusion into the obsidian, as the conventional model suggests, the depth distributions

of hydrogen and  $^{18}\text{O}$  measured by SIMS should strongly correlate or coincide. Our experimental temperature range was chosen to be from 75 to  $200^\circ\text{C}$ ; *i.e.*, high enough to accelerate the hydration process and get an adequate hydration experiment time, but far from the obsidian transition temperature of  $\sim 400^\circ\text{C}$ , above which another mechanism of the water - obsidian interaction should be considered [12].

### 2. Experimental

Experimental obsidian samples about  $3 \times 5 \times 2$  mm in size were cut by a diamond saw from a big obsidian rock found in Cerro de las Navajas, Central Mexico (see Fig. 1A). Then the samples (Fig. 1C) were polished by the standard technique; we used sequentially meshes of 800 grits, 1200 grits, and 1600 grits, and then finished with a cloth and a 1-micron diamond paste. The sample hydration was performed in standard 1.33-inch conflate nipples sealed with copper gaskets which supported heating up to  $250^\circ\text{C}$  (see Fig. 1B). The top flange was drilled, a hole was threaded, and a stainless steel screw was used to hang up the obsidian samples with a Teflon wire. Therefore, any contact of the samples with water and with the walls was avoided. Before tightening the screws, thus fixing the top flange with the samples, approximately 500 milligrams of isotopic water with 97% of the  $^{18}\text{O}$  isotope from Sigma Aldrich (p/n 329878-1G) was poured on the nipple bottom. Then the nipple with the obsidian sample was loaded vertically into an oven with a digital temperature controller pre-heated to the desired temperature where it was kept for the time from 2 days ( $200^\circ\text{C}$ ) and up to 5 months ( $90^\circ\text{C}$ ). After hydration the sample was withdrawn from the nipple, cleaned in alcohol in an ultrasonic bath during 5-10 minutes, and dried in a dry nitrogen flux. Then the sample was fixed in a special sample holder and loaded into a load-lock camera of the SIMS instrument for more than 8 hours and pumped down to  $10^{-6}$  Pa. The hydration experiments were repeated

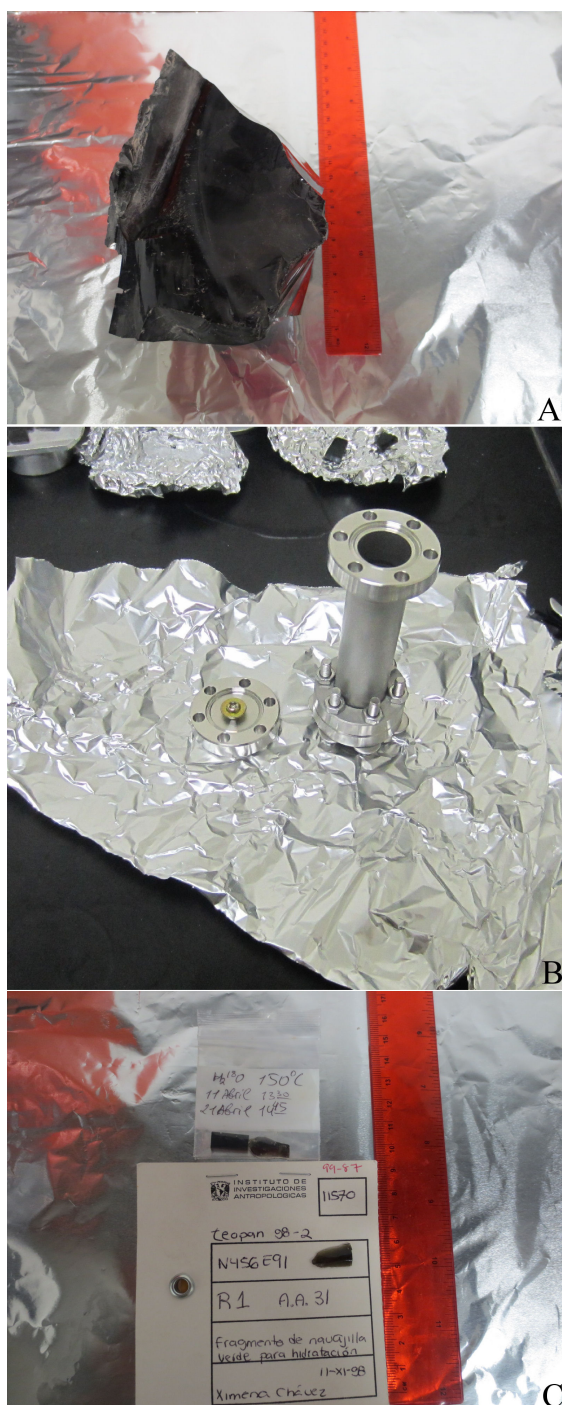


FIGURE 1. Photos of obsidian rock (A); 1.33'' Conflate nipples (B) used for sample preparation; archaeological (above) and hydrated in laboratory (below) experimental samples (C).

for each temperature to confirm a reproducibility of the method. The list of the experimental samples with the time and the temperature of hydration anyone can find in the Table I.

The chemical composition of the obsidian samples was analyzed independently by the XRF method at the Geoarchaeological XRF Laboratory, NM, USA. Table II shows concentrations of the main components detected by XRF.

TABLE I. The list of laboratory-hydrated obsidian samples.

N	Temperature, °C	Time, seconds	Dissolved layer, nm
1	75	$9.4 \times 10^6$ , $8.9 \times 10^6$	43
2	90	$6.4 \times 10^6$ , $1.2 \times 10^7$	nd
3	104	$2.4 \times 10^6$	76
4	117	$2.4 \times 10^6$	45
5	144	$1.9 \times 10^5$ , $3.3 \times 10^5$	34
6	167	$5.7 \times 10^5$	91
7	204	$2.6 \times 10^5$	35
8*	90	$4.2 \times 10^5$ , $3.6 \times 10^5$	nd

\*Soda-lime glass sample.

It is well known that the water-obsidian interaction results, among other effects, in dissolution or weathering of the obsidian surface. We examined the dissolution of obsidian during our hydration experiment by covering a part of the experimental sample surfaces with a 100 -nm gold film. In our experiments with different temperatures the dissolved layer did not exceed 91 nm at 167°C (see Table I), which was negligibly small as compared with the few-micron depth of penetration of  $^{18}\text{O}$  and  $^1\text{H}$  observed in our experiments.

We performed as well a “de-hydration” experiment. One set of obsidians was hydrated at 90°C during two months and another set was hydrated at 200°C during two days. And then the first set of obsidians was heated in vacuum during two months at 0.08 MPa, and the second one was heated during 8 hours under a high vacuum of  $10^{-6}$  Pa.

Several obsidian artifacts found in Teopancazco (Teotihuacan, Mexico) were analyzed by SIMS to compare with the laboratory hydrated samples.

The depth profiling analysis of the hydrated obsidians was carried out with a TOF-SIMS V instrument from Ion-TOF GmbH. A double ion beam irradiation regime was used for the measurements. Ion sputtering of the sample surface was performed by a 2 keV cesium ion beam at 45° with respect to the surface normal. The ion current reached 150-170 nA. The SIMS analysis was performed with a pulsed  $\text{Bi}^+$  ion beam with an ion energy of 30 keV and angle of incidence of 45° with respect to the normal incidence. The bismuth ion current was varied from 0.2 to 1.2 pA. The cesium beam sputtered a raster of  $300 \times 300$  microns; whereas the bismuth beam scanned a  $100 \times 100$  micron raster at the center of the cesium sputter crater. Secondary negative ions emitted from this  $100 \times 100$  micron area were separated by mass during their flight in a reflectron-type mass analyzer. The ion irradiation of obsidians (as well as any dielectric) resulted in the surface charging effect, which did not allow us to continue the analysis. To compensate for the arising positive surface charge, we used a flood electron gun with an electron beam current of 10-20  $\mu\text{A}$  and electron energy of 20 eV. All the measurements were performed under an ultrahigh vacuum of  $<1 \times 10^{-7}$  Pa. The hydrogen concentration was

TABLE II. Chemical composition of the obsidian rock used for laboratory hydration determined by XRF, in weight percent. Only the main oxides are listed in the Table.

	Na <sub>2</sub> O	Al <sub>2</sub> O <sub>3</sub>	SiO <sub>2</sub>	K <sub>2</sub> O	CaO	TiO <sub>2</sub>	MnO	Fe <sub>2</sub> O <sub>3</sub>
Test 1	4.414	9.877	77.201	4.632	0.227	0.140	0.090	2.983
Test 2	4.490	10.036	76.605	4.779	0.237	0.123	0.092	3.164

recalculated by using the implanted standard; the implantation energy of hydrogen ions was 30 keV, the dose was  $10^{17}$  ions/cm<sup>2</sup>. The concentration of the <sup>18</sup>O isotope was estimated from the standard oxygen isotope ratio by taking into consideration the total oxygen concentration in the obsidian obtained by the XRF method.

After hydration and SIMS depth profiling experiments we measured the experimental craters as well as gold steps by a surface profiler Dektak-XT from Bruker with an experimental error of 4%.

### 3. Experimental Results and Discussion

Figure 2A-F shows depth distributions (concentration vs. depth) of <sup>1</sup>H and <sup>18</sup>O in several obsidians hydrated at different temperatures. Figure 2G shows <sup>1</sup>H and <sup>18</sup>O depth distributions in the archaeological obsidian artifact found in Teopancazco with an estimated age of around 300 A.D.

First of all, we note radically different distributions of hydrogen and <sup>18</sup>O, which result, in our opinion, from two different physico-chemical processes. The hydrogen distribution, called the S-curve, is similar to those observed by other authors for laboratory-hydrated samples and for different archaeological samples [3,9-12]. The <sup>18</sup>O distribution is well fitted by the complementary error function presented in Fig. 2A-F by dashed lines. This function is the solution of the Fick's equation for a semi-infinite volume in the case of diffusion from a thin layer with a constant diffusion coefficient [13]:

$$C(t) = \frac{C_1}{2} \cdot \operatorname{erfc} \left( \frac{x}{2 \cdot \sqrt{(D \cdot t)}} \right) + C_0, \quad (2)$$

where  $C_1$  is the <sup>18</sup>O isotope concentration in the top surface layer saturated with the isotopic water, and  $C_0$  is the <sup>18</sup>O concentration in the obsidian volume due to an intrinsic oxygen.

Second, there is a top surface layer with a thickness of about 10 nm with high <sup>18</sup>O and hydrogen (or, probably, water) concentrations (see Fig. 2D). A slight leaching in this layer is observed (see the sodium depth distribution in Fig. 2D).

The experimental SIMS data can be interpreted if we assume that water molecules penetrate into the top surface layer alone, and then water molecules decompose into oxygen and hydrogen atoms due to chemical reactions [14,15], and (or) due to the water electric dipole interaction with local charges in obsidians reported by Morita *et al.* [16]. After this the

atomic hydrogen and <sup>18</sup>O isotope penetrate into the obsidian. The difference between the hydrogen and <sup>18</sup>O distributions increases with temperature; the top surface concentrations of <sup>1</sup>H and <sup>18</sup>O increase with temperature as well. These facts suggest that (i) the water concentration and also water decomposition at the obsidian surface grow with temperature, (ii) the two processes involved in the penetration of <sup>18</sup>O and <sup>1</sup>H atoms into the obsidian volume have different activation energies. Thus, we arrive at the conclusion that the original obsidian hydration model is incorrect. Atomic hydrogen rather than water penetrates into the obsidian and forms a hydrogenated layer which has been identified so far as a "hydrated" layer.

We repeated the hydration experiment for a soda-lime glass sample. The hydration was performed at 90°C during 2 months. Similar depth distributions of <sup>1</sup>H and <sup>18</sup>O in this sample are presented in Fig. 1 H, therefore the hydrated layer formation up to approximately 1-micron depth is observed. As demonstrated earlier [17,18], water penetrates into the soda-lime glass due to the ion exchange reaction between alkali oxides and water molecules.

The <sup>18</sup>O self-diffusion coefficients were obtained for different temperatures by approximating the experimental <sup>18</sup>O distribution by Eq. 2. Figure 3A shows a linear fit of experimental diffusion coefficients by the Arrhenius equation:

$$\log(D) = \log(D_0) - \frac{E_{act}}{2.303 \cdot R \cdot T}, \quad (3)$$

where  $T$  is the temperature,  $E_{act}$  is the activation energy;  $D_0$  is the pre-exponential coefficient; and  $R$  is the gas constant. The activation energy for the <sup>18</sup>O diffusion was found to be 56.73 kJ/mol, and the pre-exponential coefficient was  $D_0 = 8.97 \times 10^{-8}$  cm<sup>2</sup>/s.

In the case of hydrogen we obtained hydrogen diffusion coefficients from the SIMS data by using Eq. (1). The hydrogenated layer thickness was defined as the depth of the inflection point, where a minimum in the first derivative of the hydrogen profile was observed [10]. This point was observed at the depth where the hydration curve changed from being concave downward to concave upward. Figure 3B shows the Arrhenius plot for the hydrogen diffusion coefficients. The activation energy for the hydrogen diffusion was found to be 77.6 kJ/mol, *i.e.*, it was in a reasonable agreement with the earlier reported activation energies for "water diffusion" in obsidians which varied for different obsidians from 80 kJ/mol to 90 kJ/mol [8]. The pre-exponential coefficient for the hydrogen diffusion was  $D_0 = 8.78 \times 10^{-4}$  cm<sup>2</sup>/s, which was

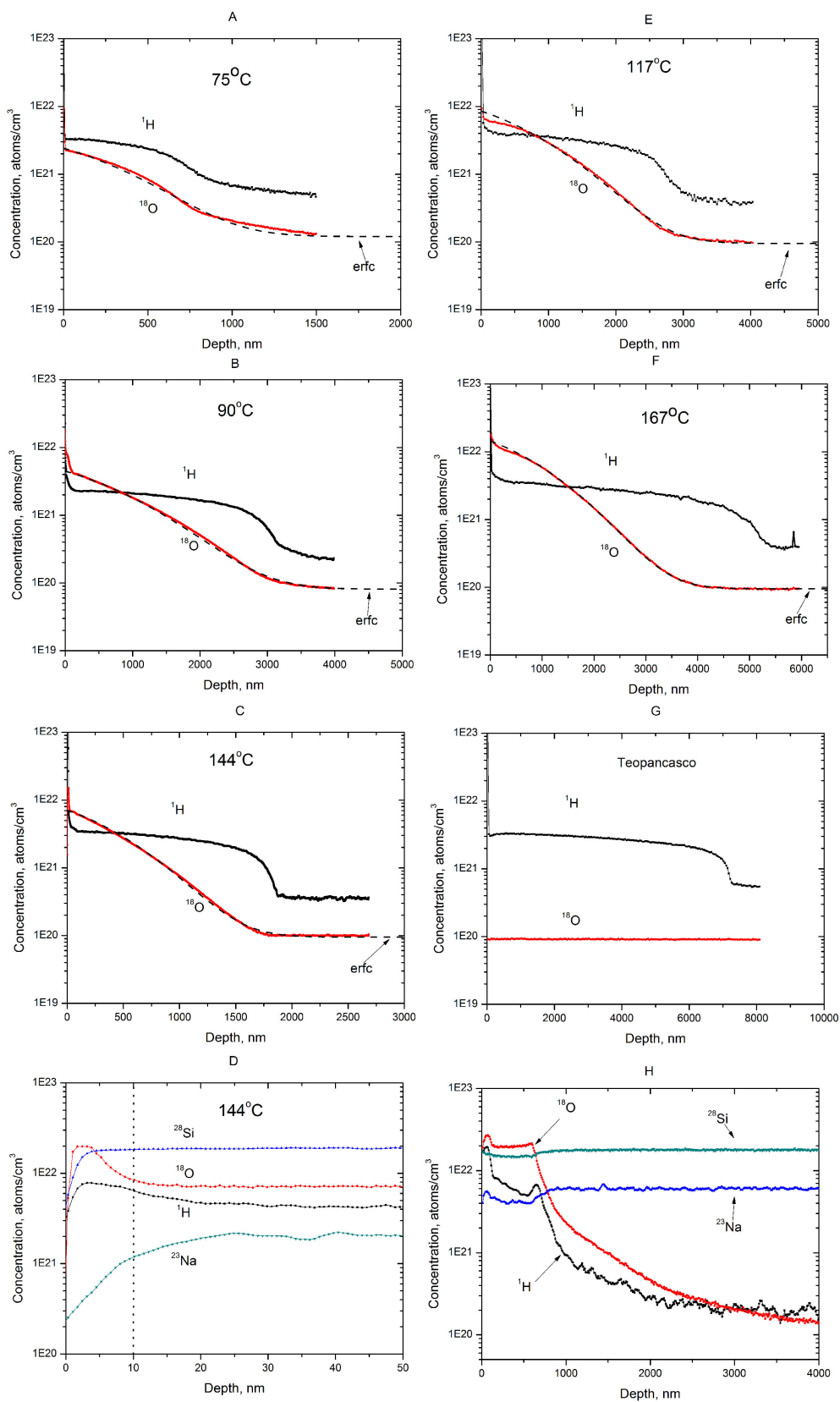


FIGURE 2. SIMS depth profiles for  $^1\text{H}$  and  $^{18}\text{O}$  in the obsidians hydrated in vapors of isotopic  $\text{H}_2^{18}\text{O}$  water at different temperatures. SIMS depth profiles for  $^1\text{H}$  and  $^{18}\text{O}$  in the obsidians hydrated in vapors of isotopic  $\text{H}_2^{18}\text{O}$  water at  $75^\circ\text{C}$  (A),  $90^\circ\text{C}$  (B),  $144^\circ\text{C}$  (C and D),  $117^\circ\text{C}$  (E), and  $167^\circ\text{C}$  (F) in comparison with the “naturally” hydrated archaeological obsidian artifact from Teopancasco (G) and with a soda-lime glass sample hydrated at  $90^\circ\text{C}$  (H) during two months.

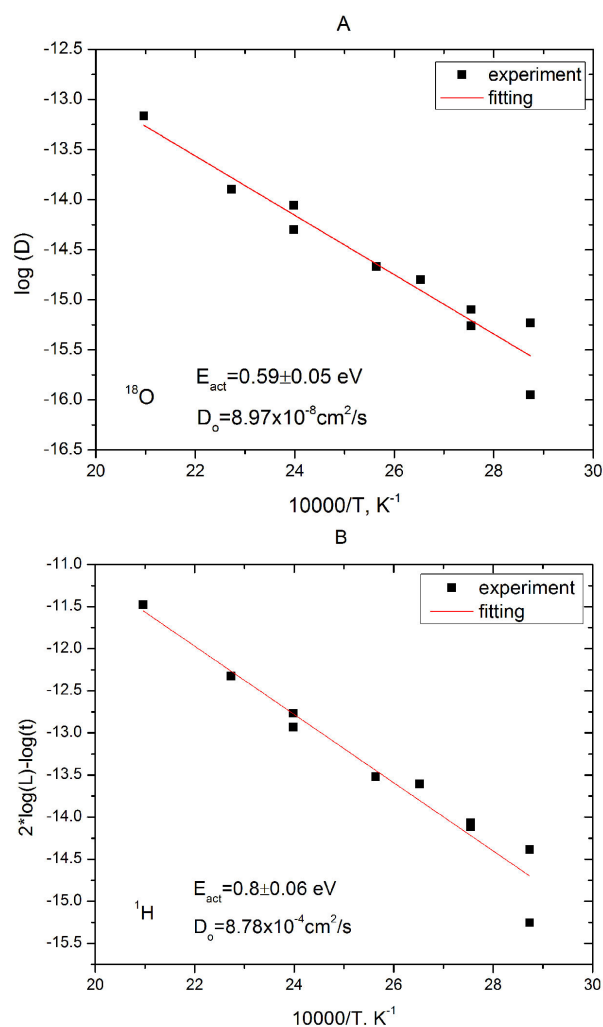


FIGURE 3. Arrhenius plot for  $^{18}\text{O}$  (A) and  $^1\text{H}$  (B) diffusion coefficients at temperatures from 75 to 204°C.

four orders of magnitude higher than for  $^{18}\text{O}$ . The obtained differences in the activation energies and the pre-exponential coefficients for hydrogen and  $^{18}\text{O}$  confirm that different physical-chemical processes occur for these elements.

Thus we established that the water-obsidian interaction results in the obsidian hydrogenation rather than hydration. But it is still unclear: what is the mechanism of hydrogen penetration into the obsidian? There are at least two possibilities: (i) a complex diffusion of hydrogen with the diffusion coefficient depending on the H concentration; and (ii) a solid state reaction between atomic hydrogen and a glass network. At the moment we have two arguments speaking in favor of a chemical reaction. The first is the absence of hydrogen effusion from hydrogenated obsidians. Riciputi *et al.* [10] mentioned that when hydrogenated obsidians were kept over 300 hours under ultrahigh vacuum ( $<10^{-6}$  Pa) at room temperature, only a slight change in the hydrogen concentration in the top surface layer of less than a 100 nm depth occurred. In our study the heating of preliminarily hydro-

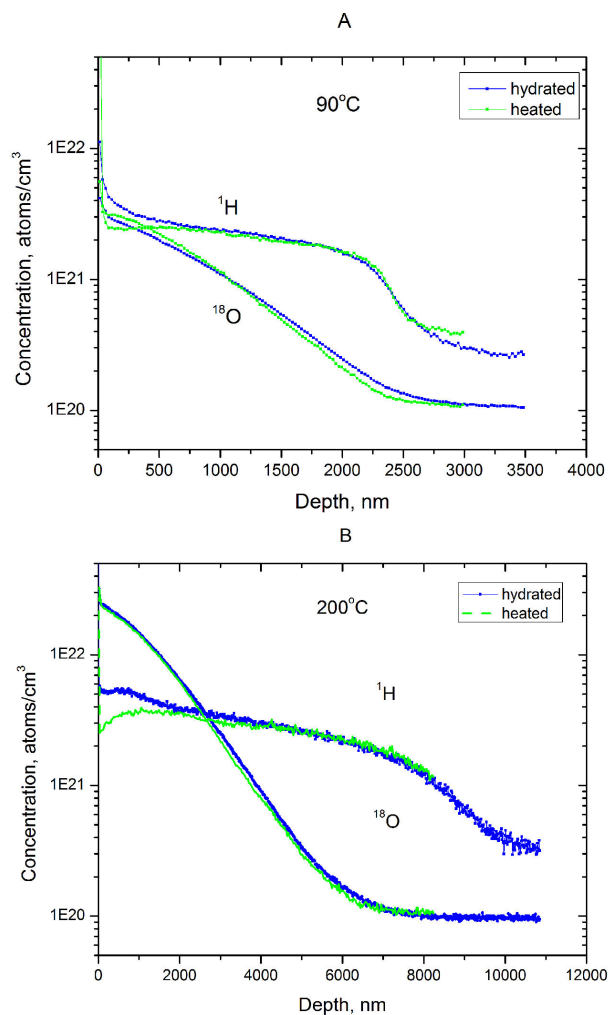


FIGURE 4. SIMS depth profiles of  $^1\text{H}$  and  $^{18}\text{O}$  in the obsidian sample hydrated in vapors of isotopic  $\text{H}_2^{18}\text{O}$  water at 90°C during two months and then heated in vacuum (0.08 MPa) at 90°C during two months (A). The same for another sample hydrated at 200°C during two days and heated at 200°C in high vacuum ( $<10^{-6}$  Pa) during 8 hours (B).

genated obsidian in vacuum during two months at 90°C and eight hours in UHV at 200°C led to a slight change in the hydrogen concentration in a top surface layer of 0.5-2  $\mu\text{m}$  (see Fig. 4 A,B).

Second, many scientists use IR spectroscopy to analyze obsidians and demonstrate a growth in the O-H and H-O-H signals in the IR spectrum with hydration time [19]. After Doremus [20], the interaction of water with silicate glasses is considered in terms of the diffusion-reaction model, which addresses water molecule interaction with silicon oxide with the formation of silicon hydroxide. In our case it could be the reaction of atomic H with those  $\text{O}_2^-$  ions of the glass network that are shared (bridging) between adjacent  $(\text{Al},\text{Si})\text{O}_4$  tetrahedra to produce OH $^-$  ions [21]. From other hand we should take into account a possible interaction of  $\text{NaAlSiO}_x$  ( $\text{KAlSiO}_x$ ) complexes with H atom resulting in formation of hydroxide of sodium and potassium [22]. We observed

a growth in the hydrogen concentration in the hydrogenated layer with temperature. In our opinion, this indicates that a change of the reaction type of hydroxides formation occurs at elevated temperatures ( $>100^{\circ}\text{C}$ ).

We have found that the water-obsidian interaction is a more complex process than that suggested earlier. During the initial period of the water-obsidian interaction the hydrogenation rate should strongly depend on the water splitting reaction and atomic hydrogen formation in the top hydrated layer of the obsidian. In this case the hydrogenated layer thickness should be linearly proportional to the time of treatment. Whereas for a long enough time the hydrogenation reaction becomes limited by the hydrogen diffusion to the hydrogenated layer - obsidian volume interface, and the growing layer thickness should be proportional to the root square of the treatment time [23]. So, the “archaeological equation” for the obsidian hydrogenation can be drawn as a linear-parabolic equation:

$$k_v \cdot L^2 + k_s \cdot L = t, \quad (4)$$

where  $k_v$  is the constant characterizing the diffusion-limited reaction of atomic hydrogen with a glass network; and  $k_s$  is the constant characterizing the surface water - glass reaction kinetics and efficiency of the atomic hydrogen formation in the top surface layer. Equation (4) is well known in the chemical community after Evans [24]; it is applied effectively, among others, to the oxide growth kinetics description. It is important to note here that the linear-parabolic equation was already suggested for obsidian “hydration” dating [25,26]. Riciputi *et al.* [10] demonstrated that the linear-parabolic equation fitted very well their experimental obsidian dating

for Chalco obsidians. This is, no doubt, an important and reassuring result from the point of view of perspectives of the new “archaeological equation” (4) for OHD.

The last but not the least remark is as follows. We observed experimentally the low-temperature water decomposition at the obsidian surface and hydrogen accumulation inside the obsidian up to a concentration of about 10 atomic %. These two experimental facts indicate that alkali-aluminosilicate glasses can be regarded as promising materials for hydrogen fusion cells due to their low cost as compared with rare-earth oxides and noble metals used at present.

## 4. Conclusions

In conclusion we would like to summarize the obtained results. We experimentally confirmed decomposition of water molecules at obsidian surface and atomic hydrogen penetration inside the obsidian with formation of the hydrogenated layer. However, contrary to the conventional model this hydrogenation layer is a result of a solid-state chemical reaction between hydrogen and glass network. We found experimentally the activation energy for the hydrogenation process, as well as for  $^{18}\text{O}$  self-diffusion in obsidian. And finally, we suggest the linear-parabolic equation as the new archaeological equation for Obsidian Hydrogenation Dating method.

## Acknowledgments

The authors thank Conacyt, Mexico for the financial support of this study, grant N 152340 and grant N 254903. The authors thank Natalie Nasina for her proofreading of the text.

1. I. Friedman, R. Smith, *American Antiquity* **25** (1960) 476-522.
2. J.R. Karsten, J.L. Delaney, *Earth Planet. Sci. Lett.* **52** (1981) 191-202.
3. L.M. Anovitz, J.M. Elam, L.R. Riciputi, D.R. Cole, *J. Arch. Sci.* **26** (1999) 735-752.
4. R. Ridings, *American Antiquity* **61** (1996) 136-148.
5. D. Webster, A. Freter, D. Rue, *Latin American Antiquity* **4** (1993) 303-324.
6. D.L. Nichols, T.H. Charlton, *Ancient Mesoamerica* **7** (1996) 231-244.
7. W.R. Ambrose, *Intrinsic hydration rate dating of obsidian*, In *Advances in Obsidian Glass Studies*, Taylor, R.E., Ed. New Jersey: Noyes Press, (1976) pp. 81-105.
8. I. Friedman, W. Long, *Science* **191** (1976) 347-352.
9. L.M. Anovitz, J.M. Elam, L.R. Riciputi, D.R. Cole, *Archaeometry* **46** (2004) 301-326.
10. L.R. Riciputi, J.M. Elam, L.M. Anovitz, D.R. Cole, *J. Arch. Sci.* **29** (2002) 1055-1075.
11. Yu. Kudriavtsev *et al.*, *Rev. Mex. Fis.* **56** (2010) 204-207.
12. L.M. Anovitz, D.R. Cole, M. Fayek, *Amer. Mineral* **93** (2008) 1166-1178.
13. J. Crank, *The Mathematics of Diffusion* Oxford: Oxford University Press, (1975), p. 32.
14. T. Yokoyama, J.F. Banfield, *Geochimica et Cosmochimica Acta* **66** (2002) 2665-2681.
15. A.F. White, *Geochimica et Cosmochimica Acta* **47** (1983) 805-815.
16. K. Morita, B. Tsuchiya, *Surf. Interface Anal* **46** (2014) 113-127.
17. W.A. Landorf, K. Davis, P. Lamarche, T. Laursen, R. Groleau, R.H. Doremus, *J. Non-Cryst. Solids* **33** (1979) 249-266.
18. M. De Bardi, H. Hutter, M. Schreiner, *Appl. Surf. Science* **282** (2013) 195-201.
19. C.M. Stevenson, S.W. Novak, *J. Arch. Sci.* **38** (2011) 1716-1726.
20. R.H. Doremus, *J. Mater. Res.* **10** (1995) 2379-2389.
21. C.W. Burnham, *Geochimica et Cosmochimica Acta* **39** (1975) 1077-1084.

22. S.C. Kohn, R. Dupree, M.E. Smith, *Geochimica et Cosmochimica Acta* **53** (1989) 2925-2935
23. N.B. Hannay, *Solid-State Chemistry*, Prentice Hall, 2-nd Edition, (1967).
24. U.R. Evans, *The Corrosion and Oxidation of Metals: Scientific Principles and Practical Applications*, London: Edward Arnold, (1960).
25. A.K. Rogers, R.M. Yohe II, *SCA Proceedings* **25** (2011) 1-15
26. F.J. Findlow, V.C. Bennett, J.E. Ericson, S.P. De Atley, *American Antiquity* **40** (1975) 344-348