

The random walk of an electrostatic field using parallel infinite charged planes

R. Aldana^a, J. Vidal Alcalá^b, and G. González^{c,d}

^a*Instituto Tecnológico y de Estudios Superiores de Occidente,
Periférico Sur Manuel Gómez Morín 8585 C.P. 45604, Tlaquepaque, Jal., México.*

^b*CIMAT y Cátedras CONACYT,*

Periférico Norte 13615, 97119 Mérida, Yucatán, México.

^c*Cátedras CONACYT, Universidad Autónoma de San Luis Potosí,
San Luis Potosí, 78000 México.*

e-mail: gabriel.gonzalez@uaslp.mx

^d*Coordinación para la Innovación y la Aplicación de la Ciencia y la Tecnología,
Universidad Autónoma de San Luis Potosí, San Luis Potosí, 78000 México.*

Received 28 October 2014; accepted 19 January 2015

We show that it is possible to generate a random walk with an electrostatic field by means of several parallel infinite charged planes in which the surface charge distribution could be either $\pm\sigma$. We formulate the problem of this stochastic process by using a rate equation for the most probable value for the electrostatic field subject to the appropriate transition probabilities according to the electrostatic boundary conditions. Our model gives rise to a stochastic law when the charge distribution is not deterministic. The probability distribution of the electrostatic field intensity, the mean value of the electrostatic force and the energy density are obtained.

Keywords: Random walk in random media; Markov chain; electrostatic field.

PACS: 41.20.Cv; O5.40.-a

1. Introduction

In a typical electrostatic problem you are given a volume charge distribution $\rho(\mathbf{r})$ and you want to find the electric field \mathbf{E} it produces. The fundamental relation between the source of the field, *i.e.* charge distribution, and the electric field is given by Gauss's law [1]

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad (1)$$

where ϵ_0 is the electric permittivity of free space. For the particular case of infinite charge sheets the electric field depends on one variable only, *i.e.* $\mathbf{E}(z)$, and the field strength it produces is constant for all distances from the plane. The electric field caused by an infinite sheet of volume charge density $\rho(z) = \sigma\delta(z)$ is given by [2]

$$\mathbf{E} = \frac{\sigma}{2\epsilon_0} \text{sign}(z) \quad (2)$$

where σ is the surface charge density and the function sign is defined as [6]

$$\text{sign}(z) = \begin{cases} +1 & \text{if } z > 0 \\ -1 & \text{if } z < 0, \end{cases} \quad (3)$$

and has been introduced to account for the vector nature of the electric field. The total electric field due to N of these infinite charge planes is then just the sum of the individual contributions, *i.e.*

$$\mathbf{E} = \sum_{n=1}^N \frac{\sigma(z_n)}{2\epsilon_0} \text{sign}(z - z_n), \quad (4)$$

where $\sigma(z_n)$ represents the surface charge density of the n -th charge plane at position z_n . Note that it is crucial to know the surface charge density on each infinite sheet to completely determine the value of the electric field through Eq. (4), if the source of the electric field is not known then we would need to approach the problem in a different way. We can formulate this same problem from a probabilistic point of view using random walks. The connection between electrostatics and Brownian motion or random walk was first shown in 1944 by Kakutani [7]. Since then many investigators have applied the probabilistic potential-theory to solve differential and integral equations for calculating electrostatic field and capacitance in integrated circuits [8, 9]. More recently, the connection between electrostatics and quantum mechanics was first shown by one of the authors [10, 11]. The purpose of this paper is to give a description of the electrostatic field when the source of the field is not known. If the volume charge density is not known then we can not use Eq. (1) to solve for the electric field, instead we need to use the boundary conditions for the electric field. For the restricted case of electrostatic charges an fields in vacuum, the appropriate boundary conditions for the electric field across a surface charge distribution is given by [1]

$$\mathbf{E}_{\text{above}} - \mathbf{E}_{\text{below}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}} \quad (5)$$

where $\hat{\mathbf{n}}$ is a unit vector perpendicular to the surface. For the case of parallel infinite charged planes the boundary condition is given by

$$E(z)_{\text{above}} - E(z)_{\text{below}} = \sigma/\epsilon_0 \quad (6)$$

Equation (6) says that, for the special case in which the infinite charged planes have a surface charge distribution given by σ , the future electric field value depends only on the present electric field value. This means that given precise information of the present state of the electric field value, the future electric field value does not depend on the past history of the process. This point is at the very heart of the Markov chain processes. This means that we can use a probabilistic model which gives a complete description of the electric field when the source of the field is not known.

The problem that we would like to address in this paper is the one where you have several parallel infinite charged planes placed along the z axis in which the surface charge density could be either $\pm\sigma$. Then the field would evolve along the z axis making random jumps each time it crosses an infinite charge sheet. We are going to show that the behavior of the electric field can be studied like a random walk. The motivation of this paper is due to the fact that many physical systems display a non-deterministic disorder. A better understanding and prediction of the nature of these systems is achieved by considering the medium to be random. The randomness models the effect of impurities on a physical system or fluctuations. Randomness appears at two levels in our problem. It comes in the description of the intensity of the electric field and also comes in the description of the medium.

The article is organized as follows. In Sec. 2 we give a brief review of Markov chains. Then in Sec. 3 we will apply the Markov chain theory to the case where we have several infinite charged planes placed along the z axis in which the surface charge density can be either $\pm\sigma$. In Sec. 4 we present the numerical results and use our model to calculate the mean electrostatic force and energy density for our problem. In the last section we summarize our conclusions.

2. Markov chains

In this section we will give a brief explanation of the Markov chain theory concepts that we need for the formulation of the problem. To introduce these concepts lets imagine a random experiment like throwing a normal dice. In this case we have six possible outcomes which are 1, 2, 3, 4, 5 or 6. We will denote the set of the possible outcomes with the letter S and name it as the *state space*. These outcomes define a random variable X which can take values from S . Each outcome can be associated with their respective probability p which must take values only in the interval $0 < p < 1$. We can write $P(X = i) = p_i$ which means that the probability of the random variable X to take a value i its equal to p_i with $i \in S$ [14].

Also if the same experiment is repeated n times, and in general each trial is dependent of the previous trials then $p_i(n)$ represents the state of the random variable X_n after n trials. If the state at a trial n is i , the probability that the next state is equal to j is defined by [3, 15]

$$P(X_{n+1} = j | X_n = i_n, X_{n-1} = i_{n-1}, \dots, X_0 = i_0) = P(X_{n+1} = j | X_n = i) = p_{ij} \quad (7)$$

Equation (7) is known as the Markov property and says that the state of a random variable after $n + 1$ transitions, only depends on the state of the random variable after n transitions. In other words, the future of the system will only depend on the present state [3, 16].

If we want to know the probabilities of a single state it is useful to use the total probability theorem which is given by [16]

$$P(B) = \sum_i P(B|A_i)P(A_i) \quad (8)$$

Where the index i runs for all the states/ A_i that B depends on. If B is the event of $X_{n+1} = j$ then $P(X_{n+1} = j) = p_j(n + 1)$. A_i is the event of $X_n = i$, thus $P(X_n = i) = p_i(n)$, then the theorem transforms into the following equation [4, 17]

$$p_j(n + 1) = \sum_i p_{ij}(n)p_i(n) \quad (9)$$

This notation allows us to interpret equation (9) as a multiplication of a horizontal vector formed with the probabilities of some certain state with a matrix formed by the p_{ij} . So if there are N possible states, $\mathbf{p}(n) = [p_{i_0}(n), \dots, p_{i_{N-1}}(n)]$ and \mathbf{W} is the matrix with elements p_{ij} , the relation between states can be written as [15, 17]

$$\mathbf{p}(n + 1) = \mathbf{p}(n)\mathbf{W}(n) \quad (10)$$

Given the initial condition $\mathbf{p}(0)$ we can determine the solution to Eq. (10) [5].

3. The Model

Suppose that we are given a collection of N infinite charged planes, half of them have a constant charge distribution σ and the other half have a constant charge distribution $-\sigma$. We can write this condition as

$$\sum_{n=0}^{N-1} \sigma(z_n) = 0 \quad (11)$$

If we randomly place the N infinite charged planes parallel to each other along the z axis at position $z_n = n \in \{0, 1, \dots, N - 1\}$, there is no way to know the values of the electric field in between the planes unless we measure the charge on the planes. Instead of that we will use Markov chain theory to analyze the problem.

We will first consider a system of four charged planes. $E(z)$ equals zero at the left and right side of the configuration due to the neutrality of the system [1]. After crossing the first charged plane, $E(z)$ would increase or decrease in an amount of σ/ϵ_0 depending if the first charged plane had a positive or negative charge, and the same for the remaining planes according to Eq. (6). In Fig. 1 there is an example of

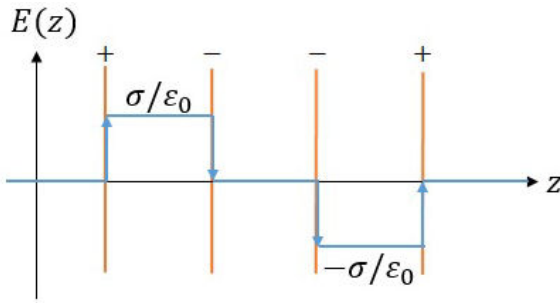


FIGURE 1. Values for $E(z)$ inside a given configuration of charged planes.

how the electric field behaves for a known charge density. In Fig. 2 we show the possible values the electric field can take when the charge density is not known.

For the sake of simplicity, we normalized the electric field dividing by σ/ϵ_0 so

$$\bar{E}(z) = \frac{\epsilon_0}{\sigma} E(z) \tag{12}$$

Due to this normalization $\bar{E}(z)$ takes only integer values i . Defining $\bar{E}_i(n)$ as our random variable, where n is the region between z_n and z_{n+1} , our state space S is defined by integer values between the interval $[-N/2, N/2]$. To calculate the transition probabilities p_{ij} we use the fact that the boundary condition in Eq. (6) prevent us to have a value bigger than $i + 1$ or smaller than $i - 1$ at a place $z = n + 1$ when $\bar{E}_i(n) = i$. This results in $p_{ij} = 0$ for $j > i + 1$ and for $j < i - 1$. After crossing a charged plane the field has to have a different value, so $p_{ii} = 0$. The last values for p_{ij} are determined by the probability of finding a negative or positive charged plane at the transition point. Defining $W_{\pm}(n)$ as the probability to find a positive or negative charged plane at a place $z_n = n$, we get the following results for p_{ij}

$$p_{ij} = \begin{cases} W_+(n) & \text{if } j = i + 1 \\ W_-(n) & \text{if } j = i - 1 \\ 0 & \text{otherwise} \end{cases} \tag{13}$$

A simple way to calculate the factors $W_{\pm}(n)$ its by dividing the remaining positive or negative planes by the remaining total planes. Let p_{\pm} be the number of positive (negative) planes between $0 \leq z \leq n$, then $p_+ + p_- = n$. Also the value of the electric field i at a point n is completely defined by the amount of positive and negative plates that have left behind, therefore $p_+ - p_- = i$. Using these two relations we can solve for p_+ :

$$p_+ = \frac{n + i}{2} \tag{14}$$

The remaining positive planes are $N/2 - p_+$ and the remaining total planes are $N - n$, then W_+ is given by

$$W_+(n) = \frac{N/2 - \frac{n+i}{2}}{N - n} = \frac{1}{2} \left(1 - \frac{i}{N - n} \right) \tag{15}$$

and $W_-(n) = 1 - W_+(n)$.

We known that the system is neutral, so the electric field must be zero at $z = 0$. In other words the probability for \bar{E} to be zero at $z = 0$ equals one, and the other cases have 0 probability. Using our notation the last statement can be written as $p_0(0) = 1$ and $p_i(0) = 0$ for $i \neq j$.

With all these information we can use Eq. (10) to obtain the most probable value of the electrostatic field in between the charged planes.

4. Results

In this section we present our results for a the system which consists of 100 charged planes ($N = 100$). Figure 3 shows a 3D plot for $\mathbf{p}(n)$. Note that the highest values for the probability are reached at $n = 0$ and $n = N - 1$ with $i = 0$, which shows the fact that the electric field must be zero at these points. In the direction of the increasing n the probabilities spreads out on the possible states, so the electric field becomes more uncertain at places near the middle of the configuration. After that point, the probabilities change to finally reach 1 at $n = N - 1$.

Its important to note that these probability functions for each transition are discontinuous functions. The plot in Fig. 4 represents a cross section for $\mathbf{p}(n)$ at the middle of the configuration ($n = 50$). In this plot it is easy to see that there are jumps between probability zero and nonzero probability along the possible states. This is due to the boundary condition given in Eq. (5). Note that the field cannot remain the same so there are forbidden states at a certain transition n .

Now we can calculate the two physical quantities that we are interested in which are the mean electrostatic force and mean electrostatic energy density. It is well known that the electrostatic force is given by $\mathbf{F} = q\mathbf{E}$ and the energy density by $\delta_E = \epsilon_0 E^2(z)/2$. In order to calculate these two quantities we need to calculate the first and second moments for the electrostatic field, *i.e.* $\langle \bar{E}(n) \rangle$ and $\langle \bar{E}^2(n) \rangle$. To calculate such moments we use the k -th moment formula [4, 16]

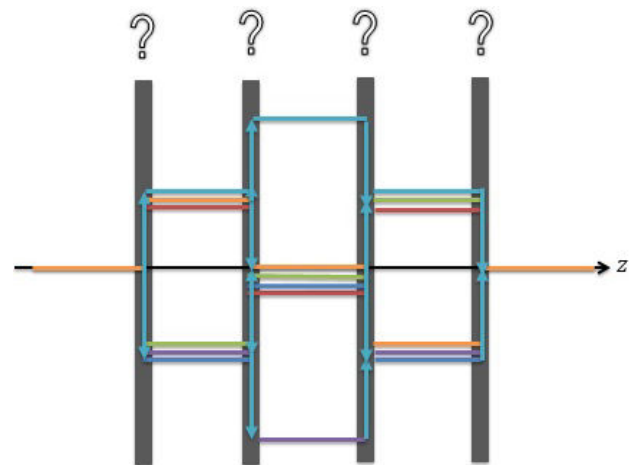


FIGURE 2. Possible values for $E(z)$ inside a four charged planes system where the charge distribution is not known.

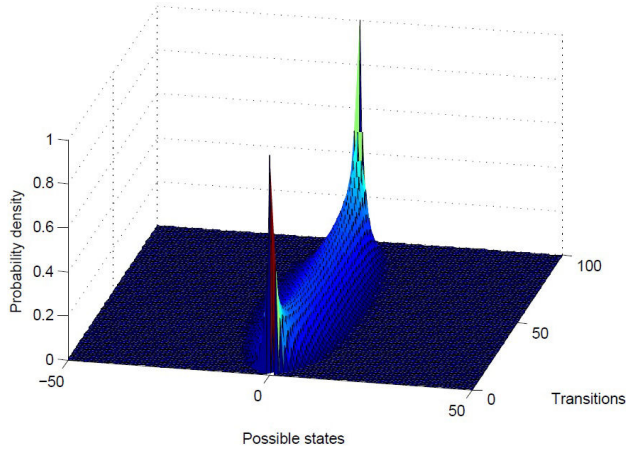


FIGURE 3. State probabilities for every transition.

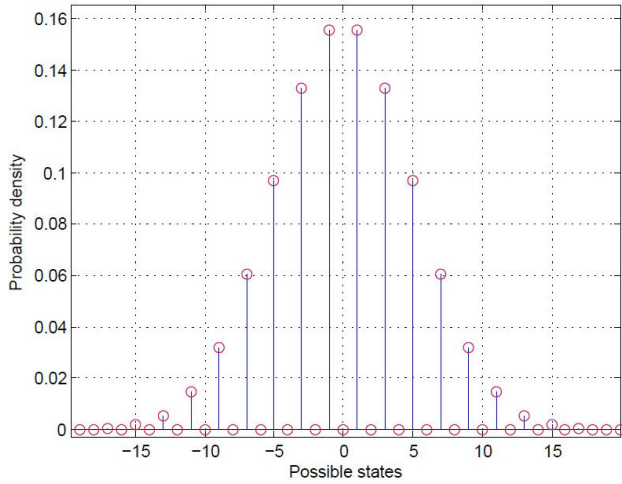


FIGURE 4. State probabilities after 50 transitions.

$$\langle \bar{E}^k(n) \rangle = \sum_i i^k p_i(n) \quad (16)$$

Where the index i runs through all S . In order to calculate these moments let us rewrite the form of Eq. (10) as follows. Most of the probabilities p_{ij} equals to zero, so the probability $p_i(n + 1)$ can be calculated adding the 2 possible paths from the last transition given by the possible values at the transition n . Each of the paths may be multiplied by the factors $W_{\pm}(n)$ [13] as is depicted in Fig. 5.

So the new form of Eq. (10) is given by

$$p_i(n + 1) = p_{i+1}(n)W_-(i + 1, n) + p_{i-1}(n)W_+(i - 1, n) \quad (17)$$

Multiplying both sides of Eq. (17) by i^k and summing from $-N/2$ to $N/2$ we may use the definition given in (16) to obtain

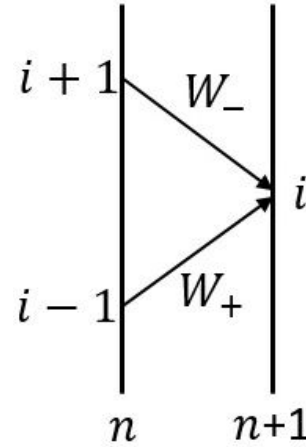


FIGURE 5. Relation between transition n and $n + 1$.

$$\begin{aligned} \langle \bar{E}^k(n + 1) \rangle &= \sum_{i=-N/2}^{N/2} i^k p_{i+1}(n)W_-(i + 1, n) \\ &+ \sum_{i=-N/2}^{N/2} i^k p_{i-1}(n)W_+(i - 1, n) \quad (18) \end{aligned}$$

Its important to note that $p_{N/2+1}(n) = p_{-N/2-1}(n) = 0$ because $N/2 + 1$ and $-N/2 - 1$ does not exist in S . Also $W_{\pm}(\pm N/2, n) = 0$ because the field cannot get a value higher (lower) than $\pm N/2$. Using these facts and setting $k = 1$ in (18) we obtain

$$\langle \bar{E}(n + 1) \rangle = \frac{N - n - 1}{N - n} \langle \bar{E}(n) \rangle \quad (19)$$

Due to the neutrality of the system we can write $\langle \bar{E}(0) \rangle = 0$. Using this initial value and substituting it in (19) its obvious that $\langle \bar{E}(n) \rangle = 0$ for all n inside the plane configuration as expected [12]. Therefore the mean force is zero everywhere.

Setting $k = 2$ in (18) and doing a little bit of algebra, Eq. (18) turns into

$$\langle \bar{E}^2(n + 1) \rangle = \frac{N - n - 2}{N - n} \langle \bar{E}^2(n) \rangle + 1 \quad (20)$$

Let $A(n)$ be the factor that multiplies $\langle \bar{E}^2(n) \rangle$ in Eq. (19), and knowing that $\langle \bar{E}^2(0) \rangle = 0$ and using recursively the difference equation in (20) [12], the states for the second moment are given by

$$\begin{aligned} \langle \bar{E}^2(1) \rangle &= 1 \\ \langle \bar{E}^2(2) \rangle &= A(1) + 1 \\ \langle \bar{E}^2(3) \rangle &= A(2)A(1) + A(2) + 1 \\ &\dots \\ \langle \bar{E}^2(n) \rangle &= 1 + \sum_{i=1}^{n-1} \prod_{j=1}^i A(n - j) \end{aligned} \quad (21)$$

Substituting $A(n)$ in equation (21), the second moment is given by

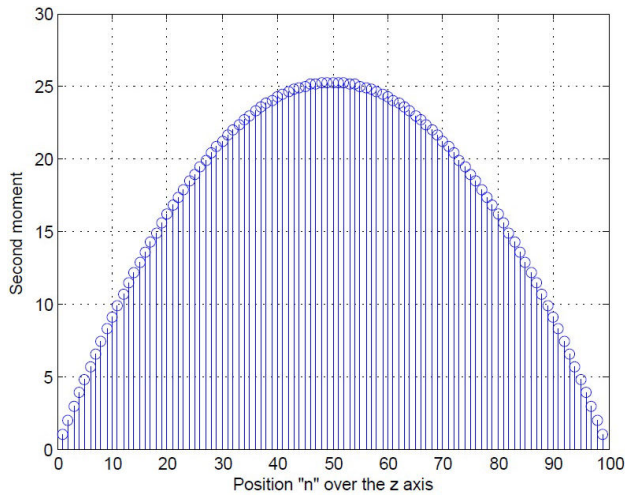


FIGURE 6. Evolution of the second moment along the z axis.

$$\langle \bar{E}^2(n) \rangle = 1 + \sum_{i=1}^{n-1} \frac{(N-n)(N-n-1)}{(N-n+i)(N-n+i-1)} \quad (22)$$

To see how the energy evolves along the z axis we used the same example with $N = 100$ as before and numerically calculated the energy using (22). The result is depicted in Fig. 6. The mean energy may be calculated using the definition of our random variable given in (12) as

$$\langle \delta_E \rangle = \frac{\epsilon_0}{2} \langle E^2(n) \rangle = \frac{\sigma^2}{2\epsilon_0} \langle \bar{E}^2(n) \rangle \quad (23)$$

We can use standard probability theory to validate our calculations for the example given above. For the case of $N = 100$ there is a probability of $50!50!/100!$ that the electrostatic field reaches the maximum (minimum) value ± 50 after 50 transitions. Using our approach this means that the probability to get a value of 50 is obtained by

$$p_{50}(50) = \prod_{j=1}^{50} W_+(50-j, 50-j) \quad (24)$$

Equation (24) means that we need 50 positive charged planes in order to reach the maximum value of the electrostatic field. Using the fact that $W_+(i, n) = (1 - i/(N - n))/2$ we get $W_+(50 - j, 50 - j) = j/(50 + j)$, inserting this into Eq. (24) we get

$$\begin{aligned} p_{50}(50) &= \prod_{j=1}^{50} \frac{j}{50+j} \\ &= \frac{(1)(2)(3) \cdots (49)(50)}{(51)(52)(53) \cdots (99)(100)} = \frac{50!50!}{(100)!} \quad (25) \end{aligned}$$

Equation (25) shows that our result is consistent with the standard probability theory.

5. Conclusions

We have shown that it is possible to perform a random walk with an electrostatic field by means of several parallel infinite charged planes in which the surface charge distribution is not explicitly known. We have worked out the special case where the charged planes have a constant surface charge distribution $\pm\sigma$ and the overall electrostatic system is neutral, *i.e.* there is the same number of positive and negative charged planes. We use Markov chain theory in order to give the most probable value for the electrostatic field in between the charged planes and use these results to obtain the mean electrostatic force and energy density. Our model gives rise to a stochastic law when the charge distribution is not deterministic.

Acknowledgments

This work was supported by the program “Cátedras CONACYT”.

1. D.J. Griffiths, *Introduction to Electrodynamics* (Prentice-Hall, Upper Saddle River, 1999)
2. M. Sadiku, *Elements of Electromagnetics* (CECSA, 2002)
3. D.T. Gillespie, *Markov Processes: An Introduction for Physical Scientists* (Gulf Professional Publishing, 1992)
4. S. Karlin, Howard M. Taylor, *A First Course in Stochastic Processes* (Academic Press, 1975)
5. J.G. Kemeny, J. Laurie Snell, *Finite Markov Chains* (Springer, 1976)
6. H.J. Weber *et al*, *Mathematical Methods for Physicists, Seventh Edition: A Comprehensive Guide* (ELSEVIER, 2012)
7. S. Kakutani, *Proc. Imp. Acad.* **20** (1944) 706-714.
8. W. Yu, K. Zhai, H. Zhuang and J. Chen, *Simulation Modelling Practice and Theory* **34** (2013) 20-26.
9. W. Yu, H. Zhuang, C. Zhang, G. Hu and Z. Liu, *IEEE Transactions on Computer-Aided Design of Int. Circ. and Syst.* **32** (2013) 353-366.
10. G. González, *Am. J. Phys.* **80** (2012) 715-719.
11. V. Rokaj, F.K. Diakonov and G. González, *Am. J. Phys.* **82** (2014) 802.
12. R.P. Grimaldi, *Discrete and Combinatorial Mathematics: An applied introduction* (Pearson, Prentice-Hall, 1998).
13. D.T. Gillespie, *American Journal of Physics* **61** (1993) 595.
14. Y.K. Leong, *American Journal of Physics* **52** (1984) 749.

15. L. Rincón, *Introducción a los procesos estocásticos* (Facultad de ciencias de la UNAM, 2012).
16. D.P. Bertsekas, J.N. Tsitsiklis, *Introduction to probability*.(Athena Scientific, Belmont, Massachusetts, 2002).
17. J. Gary Propp, D. Bruce Wilson, “*Exact Sampling with Coupled Markov Chains and Applications to Statistical Mechanics*” (1996)