Probing the effect of different exchange-correlation functionals on the optoelectronic features of chalcogenide compound Ag₂O

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The primary goal of this study is to investigate the effect of different exchange-correlation functionals on the optoelectronic and elastic properties of the Ag₂O chalcogenide compound. For the electronic structures and optical spectra, the Tran-Blaha modified Becke-Johnson approach combined with GGA and GGA+U (mBJ-GGA-PBEsol and mBJ-GGA-PBEsol+U, respectively) was used. The available theoretical and experimental data for the bandgap energy were reported to determine whether there is a correlation with our results. The electronic structure revealed that our compound is a direct semiconductor at the R-symmetry point with a bandgap of 1.22 eV, which agrees well with the experimental values for the first time. The elastic constants were also evaluated using the IRelast package, which revealed that the compound was mechanically stable. Finally, the optical response was systematically studied, and it was found that Ag₂O exhibited excellent optical efficiency.

Keywords: Chalcogenide compounds; mBJ-GGA-PBEsol+U; electronic structures; optical response; DFT calculations.

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1. Introduction

Binary silver oxide is an isostructural compound that crystallizes in the Cu₂O cuprite structure within the $Pn\tilde{3}m$ space group (224) [1,2], which is a semiconductor material. Tjeng *et al.*, have experimentally explored the electronic structures of the Ag₂O compound and reported a bandgap of 1.3 eV [3]. Silver oxide has been widely studied owing to its many important applications, such as in solar energy converters [4], battery technology [5-8], antibacterial applications [9-11], and molecular sensor technology [4], for example, as an enzyme-free glucose sensor with Cu [4,12]. It has also been reported that Ag_2O can be employed in water splitting [13], optical memory [14], photography [15], organic catalysis, such as transmetalation [16], and the oxidation of aldehydes by molecular oxygen [17]. The remarkable properties of Ag_2O which have important roles in fast-ionconducting glasses of AgI-Ag₂O-B₂O₃, AgI-Ag₂O-V₂O₃, and AgI-Ag₂O-P₂O₅ are the negative thermal expansion behavior of the cuprite Ag₂O and the structural phase transition that occurs at 35 K [18-23].

In the literature, many theoretical methods within different approximations have been employed to accurately describe the different properties of this material. F. Pei *et al.*,



FIGURE 1. The Ag₂O compound with cuprite structure.

[24], used the plane-wave-pseudopotential method as incorporated in the CASTEP code to study the structural, electronic, and optical properties of Ag₂O with cubic cuprite structure. Jeremy P. Allen et al., [25], have described the electronic structure of Ag₂O using various calculation methods. The structural, electronic, elastic, thermal, and optical properties of Ag₂O were also investigated using the Wien2k code with different approximations (LDA, GGA, Engel-Vosko-GGA, and modified-Becke-Johnson-GGA) by Haleem Din and A. H. Reshak [26]. Using the projectoraugmented-wave pseudopotentials as incorporated in the VASP code within LDA and LDA+U approximations, Naoto Umezawa et al. [27], have studied the electronic structures of Ag₂O compound. As presented previously, numerous theoretical calculations of the electronic structures have been reported for this material but, none of them has exact bandgap energy compared to the experimental ones. To resolve this discrepancy, we used a new exchange-correlation functional for the mBJ potential coupled with Hubbard correction in the framework of the DFT-theory. The main objective of this study is to provide a better description of the electronic structure and optical spectra of the chalcogenide compound Ag_2O .

2. Calculation method

In all our calculations, we used the FP-LAPW method [28,29], based upon the DFT-theory [30,31], and executed in

the Wien2k code [32]. The muffin-tin sphere radii (R_{MT}) for Ag and O are 1.96 and 1.69 (a.u), respectively. The input parameters such as $R_{MT} \times K_{max}$, G_{max} , and cut-off energy are respectively 7, 12 $Ry^{1/2}$, and -7 Ry, where K_{max} is the maximum modulus for the reciprocal lattice vector, R_{MT} is the smallest atomic sphere radius, cut-off energy is the separation energy between the valence and core states, and G_{max} defines the magnitude of the largest vector in the charge density Fourier expansion. The Monkhorst-Pack method in the first Brillouin zone (BZ) was performed using 1000 k-points. Regarding the atomic positions, the Ag₂O contains two atoms, where the Ag and O atoms occupy (0,0,0) and (1/4,1/4,1/4), respectively, as shown in Fig. 1.

GGA-PBEsol-mBJ and GGA-PBEsol-mBJ+U were used to determine the optoelectronic properties. The U values were taken as $U^1 = 5.8 \text{ eV}$ [3], and $U^2 = 11.2 \text{ eV}$ [33] for Ag-4d. The elastic properties of the Ag₂O compound were studied using the IRelast package [34,35] within the GGA-PBEsol approach.

3. Results and discussion

3.1. Equilibrium ground-state and electronic structures

The equilibrium lattice constant a_0 , bulk modulus B_0 , and its pressure derivatives, were obtained by fitting the total energy as a function of the unit-cell volume using Murnaghan's equation of state [36]. In this way, the volume optimization is carried out by minimization of the total energy versus volume. The obtained results of $a_0(\text{Å})$ and $B_0(\text{GPa})$ are given in Table I with some other available theoretical and experimental results for comparison. This table shows that the a_0 and B_0 for Ag₂O are consistent with the available theoretical and experimental results.

To explore the electronic behavior of the studied compound, we calculated the density of state DOS and electronic band-structure BS along with the principal symmetry points in the BZ for the cuprite structure (X-R-M- Γ -R) using the GGA-PBEsol-mBJ and GGA-PBEsol-mBJ+U for the exchange-correlation functional (see Fig. 2). The result of BS reveals the semiconductor nature with a direct bandgap at the R-symmetry point. The distinct features noticed within GGA-PBEsol-mBJ compared to GGA-PBEsol-mBJ+U shifts the conduction band upward and increase the bandgap.

Compound		$a_0(\text{\AA})$	$B_0(GaP)$				
	This work	Exp.	Theor.	This work	Exp.	Theor.	
Ag ₂ O	4.714	4.74 [37]	4.73 [26]	90.81	-	86.73 [26]	
		4.718 [33]	4.81 [24]			74.0 [24]	
			4.83 [37]				



FIGURE 2. The calculated band structure of Ag₂O compound using: a) mBJ-GGA-PBEsol, b) mBJ-GGA-PBEsol+ U^1 , and c) mBJ-GGA-PBEsol+ U^2 approaches.

Compound	The bandgap (eV)					
	This work	Exp.	Theor.			
	1.22	1.24 [33]	0.63 (mBJ) [26]			
Ag_2O		1.30 [3]	0.41 (EV-GGA) [26]			
		1.40 [38]	0.176 (GGA) [26]			
			0.155 (LDA) [26]			
			0.17 (LDA+U) [27]			

TABLE II. The Calculated bandgap (in eV) of Ag₂O compound.

The overall bandgap result of GGA-PBEsol-mBJ is better than that of other theoretically calculated bandgaps (LDA, GGA, EV-GGA, and LDA+U) [26,27], but significantly underestimated the bandgap from the experimental values as shown in Table II. Whereas, to resolve this discrepancy of bandgap and provide a value almost comparable to the experimental data, we used the GGA-PBEsol-mBJ+U. It should be noted that the obtained value is consistent with the experimental values for the first time. Furthermore, the conductionband-minimum CBM shifts towards higher energy compared to the GGA-PBEsol-mBJ position, but the position of the valence-band-maximum VBM remained fixed.

Total and partial DOS (TDOS and PDOS, respectively) calculated using GGA-PBEsol-mBJ+ U^2 are shown in Fig. 3. The DOS spectra are divided into four regions. The First region, situated between (-6.7 and -5.8 eV) is a result of mixed Ag-4d and O-2p states. The lower limit of this region also includes a small amount of Ag-5s states. The second one located between (-5.8 eV and -2.7 eV) is dominated principally by Ag-4d states, with a minority of O-2p states. The third region below the Fermi level (E_F) consists of mixed Ag-4d and O-2p states, with a small contribution from Ag-5s. The bottom of the conduction band is mainly composed of Ag-(5s, 5p) and O-(2s, 2p) states.



FIGURE 3. The calculated total and partial density of states for Ag_2O compound using the mBJ-GGA-PBEsol+ U^2 approach.

3.2. Elastic properties

The calculation of elastic properties is important to obtain more information about the compound, such as mechanical stability, anisotropic binding character, bonding forces, and brittle or ductile nature. For cubic Ag₂O, the number of elastic constants is reduced to only three independent elastic parameters: C_{11} , C_{12} , and C_{44} (cubic crystals). The estimated values of these constants are useful for obtaining information about mechanical behavior. The calculated elastic constants are listed in Table III. We can see clearly that the obtained

TABLE III. The Calculated elastic constants $C_{ij}(C_{11}, C_{12}, C_{44})$, Bulk modulus (B), shear modulus (G), Young's modulus (E) (in GPa), Pugh's ratio (B/G), Poisson's ratio (ν) and anisotropy factor (A) of Ag₂O compound

Compound		C_{11}	C_{12}	C_{44}	В	G	Е	B/G	А	ν
Ag_2O	This work	106.69	84.54	51.14	90.81	30.04	81.16	3.02	4.61	0.35
	Theor. [26]	132.36	70.94	56.18	86.73	44.093	113.11	1.967	1.829	0.282

 C_{ij} satisfies the three well-known Born-Huang stability criteria for the cubic system [39], which means our material is mechanically stable. It is well-known that the anisotropy factor (A) is an important mechanical parameter for technological and engineering applications and for detecting micro-cracks in materials. If the anisotropy factor is equal to unity, the material is isotropic and for any value different from unity the material is anisotropic.

The anisotropy factor value is found to be 4.61 which indicates that the compound is an anisotropic material. For more information about the compound such as stiffness, type of bounds, hardness, and ductility (brittle), we have estimated the macroscopic mechanical moduli such as Young's modulus (*E*), Shear modulus (*G*), Pugh's ratio (*B*/*G*) and Poisson's ratio (ν). For that, we have applied the Voigt-Reuss-Hill approximation [40,41].

The calculated G_V , G_R , and G are 35.11, 24.98, and 30.04 GPa, respectively. Young's modulus (*E*) provides information about the stiffness of materials. For the higher value of E = 81.16 GPa, the material shows high stiffness.

In addition, Poisson's ratio (ν) provides information about the bonding forces [42]. The central forces in the solid are trapped between 0.25 and 0.5, for the studied compound, we found $\nu = 0.35$. Due to this value, we can state that this compound is characterized by a metallic bonding. The two traditional parameters, Pugh's ratio of ductility and Poisson's ratio [43,44], are responsible for the brittle or ductile of a material. The calculated B/G and ν are 3.02 and 0.35, respectively. We can classify from these calculated values the investigated compound as ductile material. Table III displays our results of elastic constants and elastic moduli with other reported theoretical results. We can observe a significant deviation compared to our results, which is mainly due to the difference in the employed approximation.

3.3. Optical properties

To obtain more information about the electronic structure, we calculated the complex dielectric function $\epsilon(\omega)$. More descriptions of the optical calculation can be found in Ref. [45].



FIGURE 4. a) Real part of the dielectric function $\epsilon_1(\omega)$. b) Imaginary part of the dielectric function $\epsilon_2(\omega)$. c) The refractive index $n(\omega)$. d) The reflectivity $R(\omega)$. e) The optical conductivity $\sigma(\omega)$.

In this study, we computed the optical spectra of $\epsilon(\omega)$ with 50000 k-points in the IBZ. The variations in the real and imaginary parts ($\epsilon_1(\omega)$, $\epsilon_2(\omega)$) of the dielectric function in the energy range of 0-12 eV are plotted in Fig. 4a) and Fig. 4b) respectively. At zero frequency, the static dielectric constant $\epsilon_1(0)$ is 3.20. The real part of the $\epsilon_1(\omega)$ spectrum grows with increasing energy until reaching a maximum value of 5.68 at 3.73 eV. After the peak, the $\epsilon_1(\omega)$ decreases with an increase in energy, becomes negative at 5.97 eV, and then increases at higher energy. The negative values of the $\epsilon_1(\omega)$ idealize the total reflection of the incident photons [46].

It can be seen in Fig. 4b) that the imaginary part has threshold energy at 1.22 eV, which represents the optical transition between VBM and CBM. We can also notice that the maximum absorption is located in the energy range from 4.19 to 6.47 eV, the highest peak intensity of $\epsilon_2(\omega)$ is 5.86 at 5.26 eV. After this point, the $\epsilon_2(\omega)$ spectrum increases.

The refractive index $n(\omega)$ of the Ag₂O compound versus photon energy is shown in Fig. 4c), the static refractive index n(0) is 1.78, which satisfies the relation $n(0) = (\epsilon_1(0))^{1/2}$. We remark that the $n(\omega)$ spectra that reach a maximum peak value of 2.42 at 4.81 eV. Beyond the maximum value, $n(\omega)$ decreases with increasing energy.

The variation in the reflectivity $R(\omega)$ is shown in Fig. 4d), at zero frequency, the reflectivity R(0) is 8.10%, and it is obvious that $R(\omega)$ takes the maximum value (33.60%) at 6.12 eV.

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As shown in Fig. 4e), we calculate the optical conductivity $\sigma(\omega)$. In the low-energy range, we can see a drastic decrease owing to the Drude interaction mechanism of light waves with conduction electrons. The maximum value of $\sigma(\omega)$ is 4165.58 Ω^{-1} cm⁻¹ at 5.27 eV.

4. Conclusion

Herein, we have used the FP-LAPW method based on the DFT-theory to study the effect of different exchangecorrelation functionals on the optoelectronic properties of the chalcogenide compound Ag₂O. From the calculated electronic band structure, it was observed that the Ag₂O compound exhibited a direct semiconductor nature with an energy bandgap of 1.22 eV, which agrees with the experimental values for the first time. The study of mechanical stability showed that the compound was stable against elastic deformation. From the obtained elastic moduli, we conclude that our compound is an anisotropic ductile material. The $\epsilon_1(\omega)$ and $\epsilon_2(\omega)$ parts of dielectric functions were estimated to identify optical transitions. $n(\omega)$, $R(\omega)$, and $\sigma(\omega)$ were also obtained and analyzed up to 12 eV. Finally, Ag₂O was found to be useful for optoelectronic applications.

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