

Optical soliton solutions of the Ginzburg-Landau equation with conformable derivative and Kerr law nonlinearity

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By using the generalized exponential rational function method, we obtain new periodic and hyperbolic soliton solutions for the conformable Ginzburg-Landau equation with the Kerr law nonlinearity. The conformable derivative is considered to obtain the exact solutions under constraint conditions. Numerical simulations are performed to confirm the efficiency of the proposed method.

Keywords: Soliton solutions; generalized exponential rational function method; Ginzburg Landau equation; conformable time fractional derivative; Kerr law nonlinearity.

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1. Introduction

The Ginzburg-Landau equation describes the optical soliton propagation through a wide range of waveguides such as crystals, optical metamaterials, optical fibers, and optical couplers [1-9]. Many powerful methods have been established to find soliton solutions of the Ginzburg-Landau equation including the modified simple equation method [10,11], the semi-inverse variational principle [12,13], the extended Jacobi elliptic function expansion method [14,15], the exponential rational function [16], the generalized exponential rational function method (GERFM) [17], among others.

Due to the complex nature of the optical soliton propagation, several works consider the fractional calculus to construct new optical soliton solutions [18-21]. Nevertheless, fractional derivatives do not obey some basic properties of integer derivative such as product rule and chain rule. Recently, a local derivative called conformable derivative has been formulated by Khalil in [22]. The conformable calculus satisfies all the properties of the standard calculus, for instance, the chain rule. This derivative can be considered to be a natural extension of the classical derivative [23-30].

Let $f : [0, \infty) \rightarrow \mathbb{R}$, the conformable derivative of a function $f(t)$ of order α , is defined as [22]

$$D_t^\alpha f(t) = \lim_{\epsilon \rightarrow 0} \frac{f(t + \epsilon t^{1-\alpha}) - f(t)}{\epsilon},$$

$$\alpha \in (0, 1], \quad t > 0. \quad (1)$$

In this paper, the conformable GERFM is employed to study the complex time Ginzburg-Landau equation with Kerr law nonlinearity [31].

$$iD_t^\alpha q + aq_{xx} + b|q|^2 q = \frac{\beta \left[2|q|^2 (|q|^2)_{xx} - [(|q|^2)_x]^2 \right] + \gamma q}{|q|^2 q^*}, \quad (2)$$

where D_t^α is the conformable derivative of order $\alpha \in (0, 1]$, x represents the non-dimensional distance across the fiber, a represents the coefficient of group velocity dispersion, and b represents the coefficient of nonlinearity. The term β arises from the perturbation effects and γ is related to the detuning effect. All above-mentioned parameters are constants real values.

2. Overview of the generalized exponential rational function method

Let us state the main steps of GERFM as follows [17]

1. Let us take into account the following nonlinear partial differential equation

$$\mathcal{L}(\Upsilon, \Upsilon_x, \Upsilon_t, \Upsilon_{xx}, \dots) = 0. \quad (3)$$

Using the transformations $\Upsilon = \Upsilon(\chi)$ and $\chi = \sigma x - \varphi t$, in nonlinear partial differential equation (3), we define attain

$$\mathcal{L}(\Upsilon, \Upsilon', \Upsilon'', \dots) = 0, \quad (4)$$

which is proposed as an ordinary differential equation; where the values of σ and φ will be determined later.

2. Consider Eq. (4) has the solution of the form

$$\Upsilon(\chi) = A_0 + \sum_{k=1}^M A_k \Psi(\chi)^k + \sum_{k=1}^M B_k \Psi(\chi)^{-k}, \quad (5)$$

where

$$\Psi(\chi) = \frac{p_1 e^{q_1 \chi} + p_2 e^{q_2 \chi}}{p_3 e^{q_3 \chi} + p_4 e^{q_4 \chi}}. \tag{6}$$

The values of constants $p_i, q_i (1 \leq i \leq 4), A_0, A_k$ and $B_k (1 \leq k \leq M)$ are determined, in such a way that solution (5) always satisfy Eq. (4). By considering the homogenous balance principle, the value of M is determined.

3. Placing Eq. (5) into Eq. (4), we give the following algebraic equation $\Xi(\Lambda_1, \Lambda_2, \Lambda_3, \Lambda_4) = 0$, in terms of $\Lambda_i = e^{q_i \chi}$ for $i = 1, \dots, 4$. After setting each of the coefficients of variables in Ξ to zero, a system of non-linear equations in terms of $p_i, q_i (1 \leq i \leq 4)$, and $\sigma, \varphi, A_0, A_k$ and $B_k (1 \leq k \leq M)$ are generated.
4. By solving the above equations systems using any symbolic computation software, the values of $p_i, q_i (1 \leq i \leq 4), A_0, A_k$, and $B_k (1 \leq k \leq M)$ are determined, replacing these values into Eq. (5) provides us the soliton solutions of Eq. (3).

3. Application

In order to find solutions of Eq. (2), the following new variables are introduced

$$\begin{aligned} q(x, t) &= \Theta(\chi) e^{i\Phi(x, t)}, & \xi &= x - \left(\frac{\nu}{\alpha}\right) t^\alpha, \\ \Phi(x, t) &= -kx + \left(\frac{w}{\alpha}\right) t^\alpha, \end{aligned} \tag{7}$$

where ν and k are the speed and frequency of the soliton, respectively; w represents the wave number of the soliton. Considering Eq. (7), we convert Eq. (2) in the following expression

$$(a - 4\beta) \Theta'' - (w + ak^2 + \gamma) \Theta + b \Theta^3 = 0, \tag{8}$$

from real part

$$\nu = -2ak, \tag{9}$$

and Eq. (9), from the imaginary parts.

If we apply the balance principle on the terms Θ^3 and Θ'' in Eq. (8), we have $3M = M + 2$, so $M = 1$. Using Eq. (6) together with $M = 1$, we have

$$\Theta(\chi) = A_0 + A_1 \Psi(\chi) + \frac{B_1}{\Psi(\chi)}. \tag{10}$$

Proceeding as outlined in second section, we obtain the following sets of solutions

Set 1:

One obtains $p = [-1, 0, 1, 1]$ and $q = [0, 0, 1, 1]$, so Eq. (6) turns to

$$\Psi(\chi) = -\frac{1}{1 + e^\chi}. \tag{11}$$

We also obtain

$$\begin{aligned} k &= k, & w &= -k^2 a - a/2 + 2\beta - \gamma, \\ A_0 &= \frac{-\sqrt{2}\sqrt{4\beta - a}}{2\sqrt{b}}, \\ A_1 &= \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, & B_1 &= 0. \end{aligned}$$

Placing values in Eqs. (10) and (11), yields the following solution

$$\Theta(\chi) = \frac{(1 - e^\chi) \sqrt{4\beta - a} \sqrt{2}}{\sqrt{b} (2 + 2e^\chi)},$$

and

$$\begin{aligned} q_1(x, t) &= \left(\frac{\left(1 - e^{\left(x + \frac{2ak}{\alpha} t^\alpha\right)}\right) \sqrt{4\beta - a} \sqrt{2}}{\sqrt{b} \left(2 + 2e^{\left(x + \frac{2ak}{\alpha} t^\alpha\right)}\right)} \right) \\ &\times e^{i\left(-kx + \left(\frac{w}{\alpha}\right) t^\alpha\right)}. \end{aligned} \tag{12}$$

Set 2:

One obtains $p = [-3, -2, 1, 1]$ and $q = [1, 0, 1, 0]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{-3 - 2e^\chi}{1 + e^\chi}. \tag{13}$$

We also obtain

$$\begin{aligned} k &= \frac{\sqrt{2}\sqrt{4\beta - a} - 2\gamma - 2w}{2\sqrt{a}}, & w &= w, \\ A_0 &= \frac{-5\sqrt{2}\sqrt{4\beta - a}}{2\sqrt{b}}, & A_1 &= 0, \\ B_1 &= \frac{-6\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}. \end{aligned}$$

Placing values in Eqs. (10) and (13), yields the following solution

$$\Theta(\chi) = \frac{\sqrt{2}\sqrt{4\beta - a} (2e^\chi - 3)}{\sqrt{b} (6 + 4e^\chi)},$$

and

$$\begin{aligned} q_2(x, t) &= \left(\frac{\sqrt{2}\sqrt{4\beta - a} \left(2e^{\left(x + \frac{2ak}{\alpha} t^\alpha\right)} - 3\right)}{\sqrt{b} \left(6 + 4e^{\left(x + \frac{2ak}{\alpha} t^\alpha\right)}\right)} \right) \\ &\times e^{i\left(-kx + \left(\frac{w}{\alpha}\right) t^\alpha\right)}. \end{aligned} \tag{14}$$

Set 3: One obtains $p = [2, 0, 1, 1]$ and $q = [-1, 0, 1, -1]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{\cosh(\chi) - \sinh(\chi)}{\cosh(\chi)}. \tag{15}$$

We also obtain

$$k = \frac{\sqrt{8\beta - 2a - \gamma - w}}{\sqrt{a}}, \quad w = w,$$

$$A_0 = \frac{\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}},$$

$$A_1 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad B_1 = 0. \quad (16)$$

Placing values in Eqs. (10) and (15), yields the following solution

$$\Theta(\chi) = \frac{\sinh(\chi) \sqrt{2}\sqrt{4\beta - a}}{\sqrt{b} \cosh(\chi)},$$

and

$$q_3(x, t) = \left(\frac{\sinh\left(x + \frac{2ak}{\alpha}t^\alpha\right) \sqrt{2}\sqrt{4\beta - a}}{\sqrt{b} \cosh\left(x + \frac{2ak}{\alpha}t^\alpha\right)} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (17)$$

Set 4:

One obtains $p = [-3, -1, 1, 1]$ and $q = [1, -1, 1, -1]$, so Eq. (6) turns to

$$\Psi(\chi) = -\frac{\sinh(\chi) + 2 \cosh(\chi)}{\cosh(\chi)}. \quad (18)$$

We also obtain

$$k = \frac{\sqrt{8\beta - 2a - \gamma - w}}{\sqrt{a}}, \quad w = w,$$

$$A_0 = \frac{2\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}},$$

$$A_1 = 0, \quad B_1 = \frac{3\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

Placing values in Eqs. (10) and (18), yields the following solution

$$\Theta(\chi) = \frac{\sqrt{2}\sqrt{4\beta - a} (2 \sinh(\chi) + \cosh(\chi))}{\sqrt{b} (2 \cosh(\chi) + \sinh(\chi))},$$

and

$$q_4(x, t) = \left(\frac{\sqrt{2}\sqrt{4\beta - a} (2 \sinh\left(x + \frac{2ak}{\alpha}t^\alpha\right) + \cosh\left(x + \frac{2ak}{\alpha}t^\alpha\right))}{\sqrt{b} (2 \cosh\left(x + \frac{2ak}{\alpha}t^\alpha\right) + \sinh\left(x + \frac{2ak}{\alpha}t^\alpha\right))} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (19)$$

Set 5:

One obtains $p = [1 - i, -1 - i, -1, 1]$ and $q = [i, -i, i, -i]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{\cos(\chi) - \sin(\chi)}{\sin(\chi)}. \quad (20)$$

We also obtain

$$k = k, \quad w = -k^2a + 2a - 8\beta - \gamma, \quad A_0 = \frac{\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad A_1 = 0, \quad B_1 = \frac{2\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

Placing values in Eqs. (10) and (20), yields the following solution

$$\Theta(\chi) = \frac{(\sin(\chi) + \cos(\chi)) \sqrt{2}\sqrt{4\beta - a}}{\sqrt{b} (\cos(\chi) - \sin(\chi))},$$

and

$$q_5(x, t) = \left(\frac{(\sin\left(x + \frac{2ak}{\alpha}t^\alpha\right) + \cos\left(x + \frac{2ak}{\alpha}t^\alpha\right)) \sqrt{2}\sqrt{4\beta - a}}{\sqrt{b} (\cos\left(x + \frac{2ak}{\alpha}t^\alpha\right) - \sin\left(x + \frac{2ak}{\alpha}t^\alpha\right))} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (21)$$

Set 6:

One obtains $p = [-2 - i, 2 - i, -1, 1]$ and $q = [i, -i, i, -i]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{2 \sin(\chi) + \cos(\chi)}{\sin(\chi)}. \quad (22)$$

We also obtain

$$k = k, \quad w = -k^2a + 2a - 8\beta - \gamma, \quad A_0 = \frac{-2\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad A_1 = 0, \quad B_1 = \frac{5\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

Placing values in Eqs. (10) and (22), yields the following solution

$$\Theta(\chi) = \frac{\sqrt{2}\sqrt{4\beta - a}(\sin(\chi) - 2\cos(\chi))}{\sqrt{b}(\cos(\chi) + 2\sin(\chi))},$$

and

$$q_6(x, t) = \left(\frac{\sqrt{2}\sqrt{4\beta - a}(\sin(x + \frac{2ak}{\alpha}t^\alpha) - 2\cos(x + \frac{2ak}{\alpha}t^\alpha))}{\sqrt{b}(\cos(x + \frac{2ak}{\alpha}t^\alpha) + 2\sin(x + \frac{2ak}{\alpha}t^\alpha))} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (23)$$

Set 7:

One obtains $p = [2 - i, -2 - i, -1, 1]$ and $q = [i, -i, i, -i]$, so Eq. (6) turns to

$$\Psi(\chi) = -\frac{2\sin(\chi) - \cos(\chi)}{\sin(\chi)}. \quad (24)$$

We also obtain

$$k = k, \quad w = -k^2a + 2a - 8\beta - \gamma, \quad A_0 = \frac{2\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad A_1 = 0, \quad B_1 = \frac{5\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

Placing values in Eqs. (10) and (24), yields the following solution

$$\Theta(\chi) = \frac{-\sqrt{2}\sqrt{4\beta - a}(\sin(\chi) + 2\cos(\chi))}{\sqrt{b}(-\cos(\chi) + 2\sin(\chi))},$$

and

$$q_7(x, t) = \left(\frac{-\sqrt{2}\sqrt{4\beta - a}(\sin(x + \frac{2ak}{\alpha}t^\alpha) + 2\cos(x + \frac{2ak}{\alpha}t^\alpha))}{\sqrt{b}(-\cos(x + \frac{2ak}{\alpha}t^\alpha) + 2\sin(x + \frac{2ak}{\alpha}t^\alpha))} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (25)$$

Set 8:

One obtains $p = [2, 0, 1, -1]$ and $q = [1, 0, 1, -1]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{\cosh(\chi) + \sinh(\chi)}{\sinh(\chi)}. \quad (26)$$

We also obtain

$$k = k, \quad w = -k^2a - 2a + 8\beta - \gamma, \quad A_0 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad A_1 = \frac{\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad B_1 = 0.$$

Placing values in Eqs. (10) and (26), yields the following solution

$$\Theta(\chi) = \frac{\cosh(\chi)\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}\sinh(\chi)},$$

and

$$q_8(x, t) = \left(\frac{\cosh(x + \frac{2ak}{\alpha}t^\alpha)\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}\sinh(x + \frac{2ak}{\alpha}t^\alpha)} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (27)$$

Figures 1(a-d) show numerical simulations of Eq. (27) for $\alpha = 1, 0.7, 0.5, 0.3$, arbitrarily chosen.

Set 9:

One obtains $p = [-1, 3, 1, -1]$ and $q = [1, -1, 1, -1]$, so Eq. (6) turns to

$$\Psi(\chi) = \frac{\cosh(\chi) - 2\sinh(\chi)}{\sinh(\chi)}. \quad (28)$$

We also obtain

$$k = k, \quad w = -k^2a - 2a + 8\beta - \gamma, \quad A_0 = \frac{-2\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad A_1 = 0, \quad B_1 = \frac{-3\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

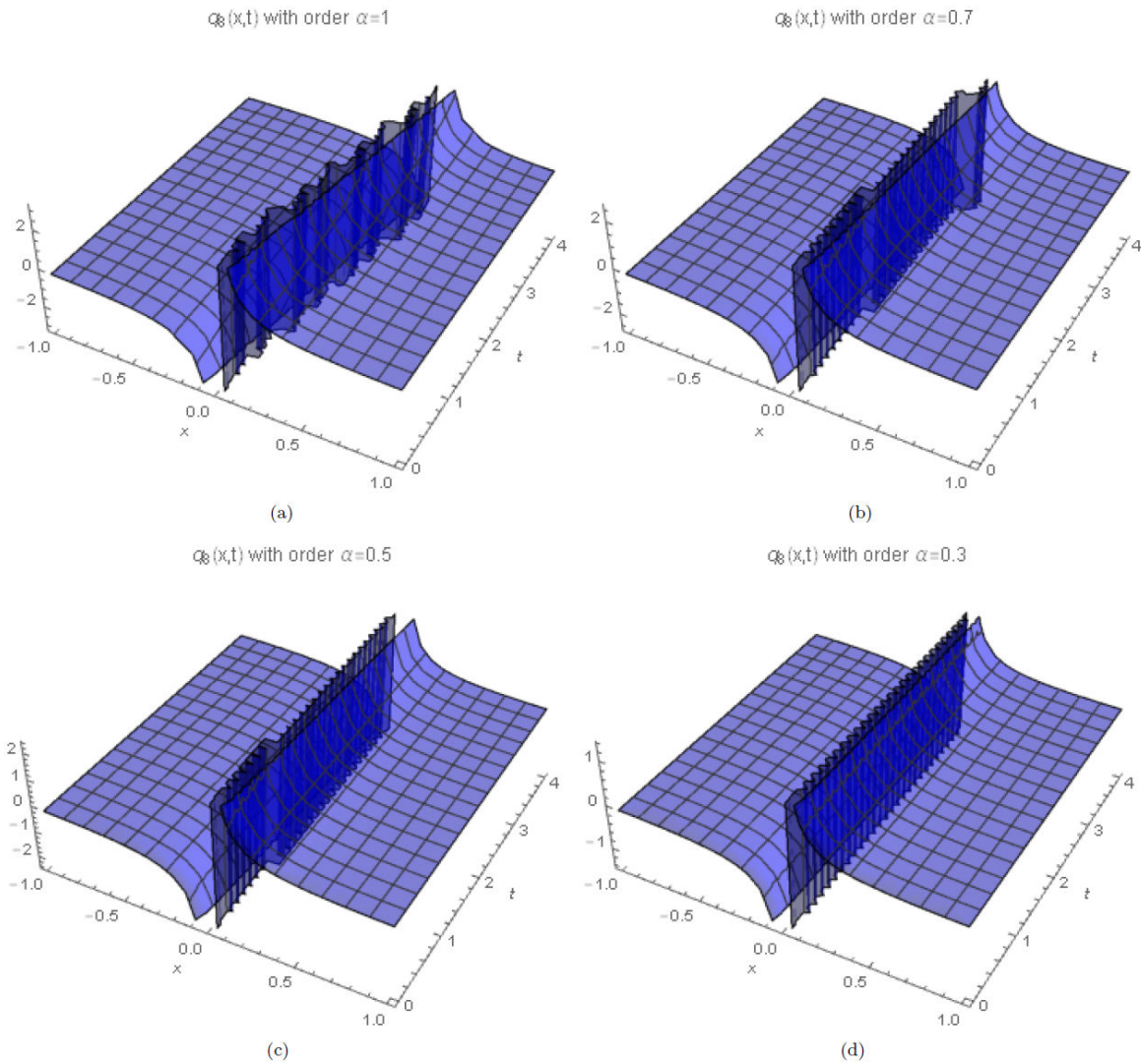


FIGURE 1. 3D Plot soliton solution related to Eq. (29).

Placing values in Eqs. (10) and (28), yields the following solution

$$\Theta(\chi) = \frac{-\sqrt{2}\sqrt{4\beta - a}(-\sinh(\chi) + 2\cosh(\chi))}{\sqrt{b}(\cosh(\chi) - 2\sinh(\chi))},$$

and

$$q_9(x, t) = \left(\frac{-\sqrt{2}\sqrt{4\beta - a}(-\sinh(x + \frac{2ak}{\alpha}t^\alpha) + 2\cosh(x + \frac{2ak}{\alpha}t^\alpha))}{\sqrt{b}(\cosh(x + \frac{2ak}{\alpha}t^\alpha) - 2\sinh(x + \frac{2ak}{\alpha}t^\alpha))} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (29)$$

Figures 2(a-d) show numerical simulations of Eq. (29) for $\alpha = 1, 0.7, 0.5, 0.3$, arbitrarily chosen.

Set 10:

One obtains $p = [i, -i, 1, 1]$ and $q = [i, -i, i, -i]$, so Eq. (6) turns to

$$\Psi(\chi) = -\frac{\sin(\chi)}{\cos(\chi)}. \quad (30)$$

We also obtain

$$k = k, \quad w = -k^2a - 4a + 16\beta - \gamma, \quad A_0 = 0, \quad A_1 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad B_1 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

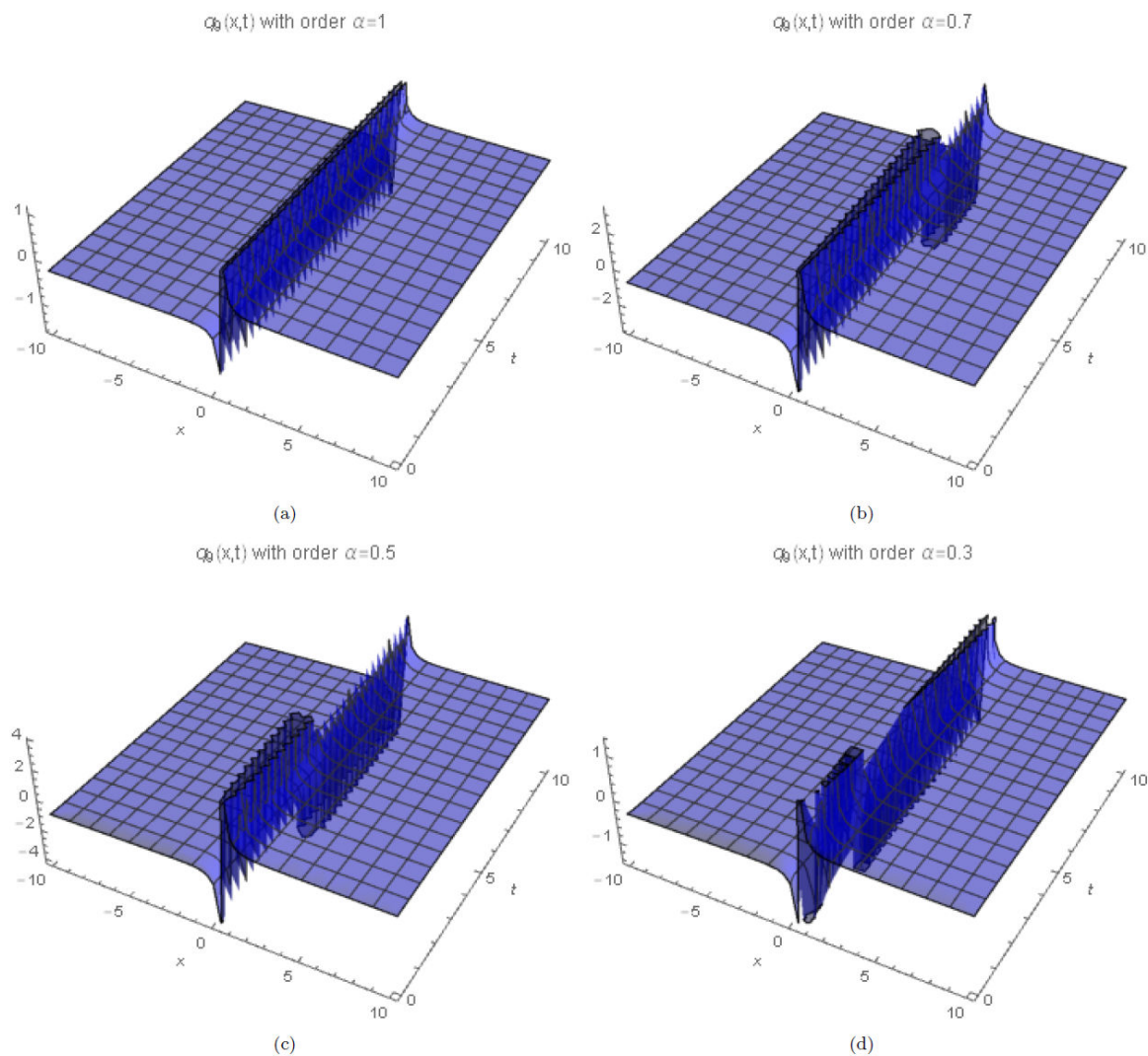


FIGURE 2. 3D Plot soliton solution related to Eq. (29).

Placing values in Eqs. (10) and (30), yields the following solution

$$\Theta(\chi) = \frac{\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}\cos(\chi)\sin(\chi)},$$

and

$$q_{10}(x, t) = \left(\frac{\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}\cos\left(x + \frac{2ak}{\alpha}t^\alpha\right)\sin\left(x + \frac{2ak}{\alpha}t^\alpha\right)} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (31)$$

Figures 3(a-d) show numerical simulations of Eq. (31) for $\alpha = 1, 0.7, 0.5, 0.3$, arbitrarily chosen.

Set 11:

One obtains $p = [1, 1, -1, 1]$ and $q = [1, -1, 1, -1]$, so Eq. (6) turns to

$$\Psi(\chi) = -\frac{\cosh(\chi)}{\sinh(\chi)}. \quad (32)$$

We also obtain

$$k = \frac{\sqrt{32\beta - 8a - \gamma - w}}{\sqrt{a}}, \quad w = w, \quad A_0 = 0, \quad A_1 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}, \quad B_1 = \frac{-\sqrt{2}\sqrt{4\beta - a}}{\sqrt{b}}.$$

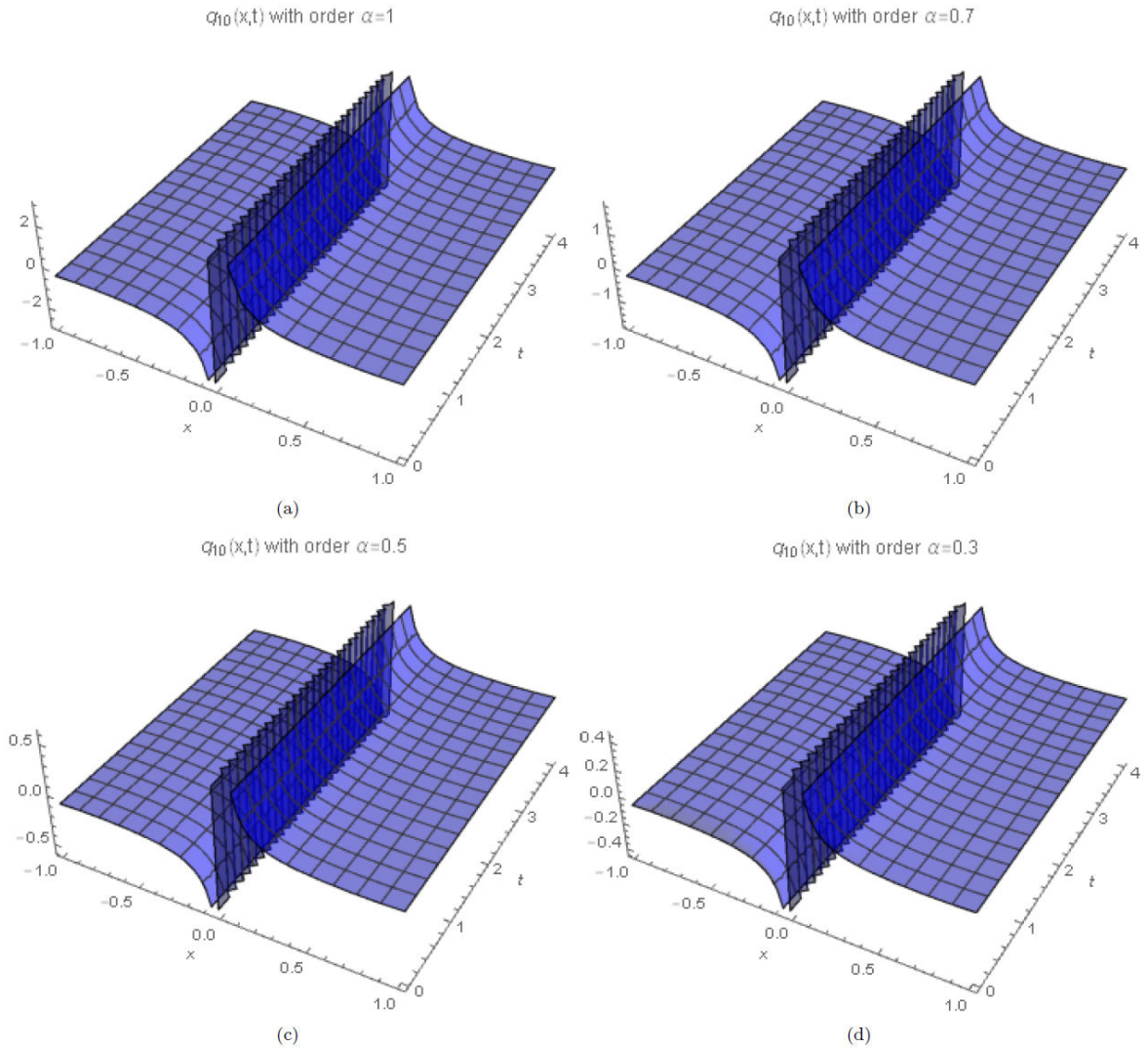


FIGURE 3. 3D Plot soliton solution related to Eq. (31).

Placing values in Eqs. (10) and (32), yields the following solution

$$\Theta(\chi) = \frac{\sqrt{2}\sqrt{4\beta - a} (\coth^2(\chi) + 1)}{\sqrt{b} \coth(\chi)},$$

and

$$q_{11}(x, t) = \left(\frac{\sqrt{2}\sqrt{4\beta - a} (\coth^2(x + \frac{2ak}{\alpha}t^\alpha) + 1)}{\sqrt{b} \coth(x + \frac{2ak}{\alpha}t^\alpha)} \right) \times e^{i(-kx + (\frac{w}{\alpha})t^\alpha)}. \quad (33)$$

Figures 4(a-d) show numerical simulations of Eq. (33) for $\alpha = 1, 0.7, 0.5, 0.3$, arbitrarily chosen.

4. Numerical simulations

In this work, we obtained different numerical solutions considering different alpha orders to obtain soliton solutions. The numerical solutions show that the change of fractional order modify the nature of the solution, and has a huge influence on the nonlinear propagation of the solitons. The analytical solutions allow graphing soliton solutions of type dark, bright,

singular or combinations. Results showed that when the time derivative decreases, the amplitude of the solitons also decreases. It happens due to the decrease in velocity of the soliton, the order α characterizing the existence of the fractional structures on the system. The new analytical soliton solutions obtained in this paper have not been reported in the literature so far. Classical soliton solutions are recovered in the limit when $\alpha \rightarrow 1$.

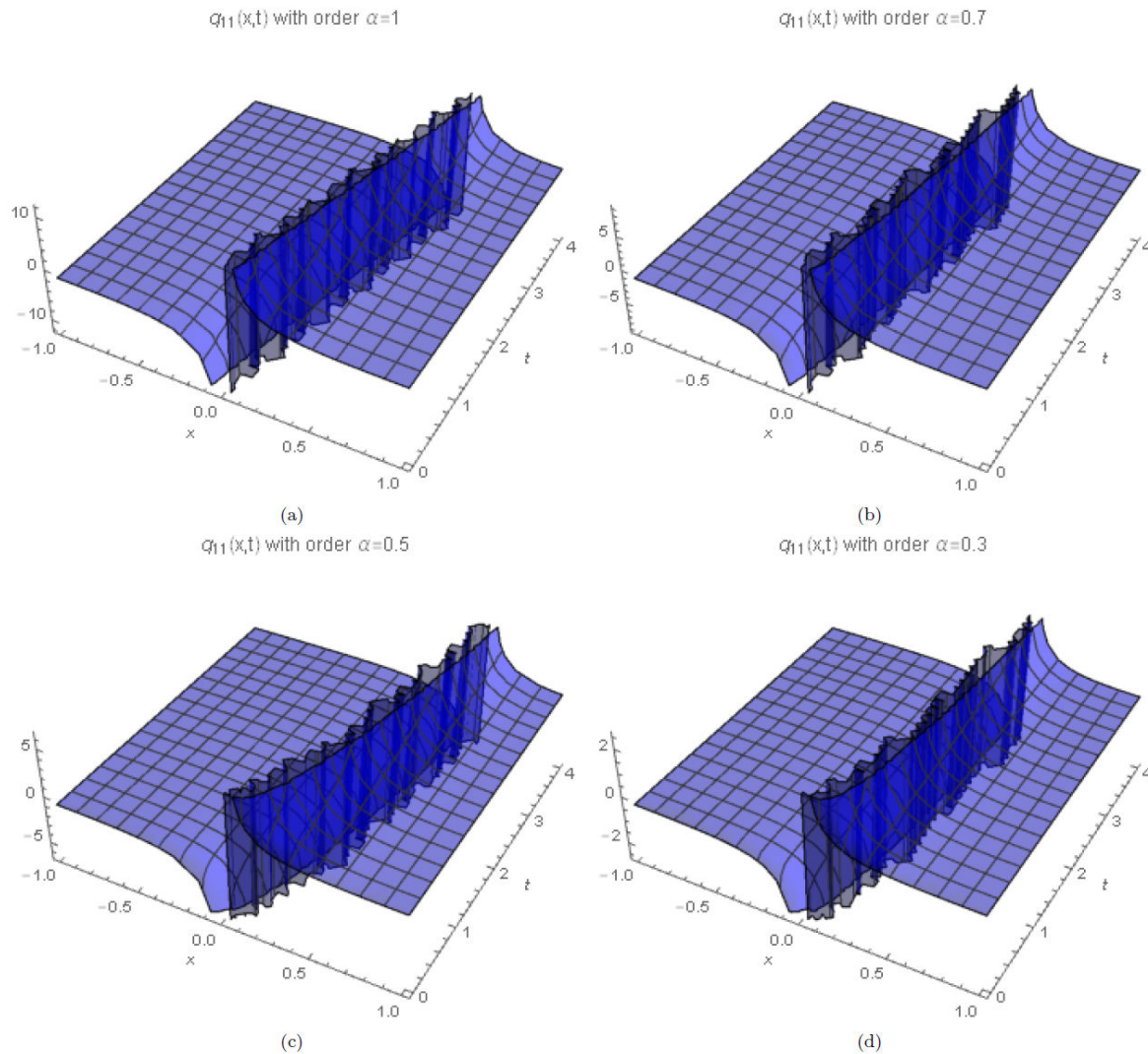


FIGURE 4. 3D Plot soliton solution related to Eq. (33).

5. Conclusion

In this work, we consider the generalized exponential rational function method to obtain approximate soliton solutions of the conformable Ginzburg-Landau equation with Kerr law nonlinearity. The Ginzburg-Landau equation describes the optical soliton propagation through a wide range of waveguides such as crystals, optical metamaterials, optical fibers and optical couplers. These soliton play an important and key role for information transfer via optical fibers.

The results showed that the generalized exponential rational function method is a promising approach to integrate the Ginzburg-Landau equation. We believe that this method also can be extrapolated to other nonlinear problems which arise in the theory of solitons.

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Competing interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

Authors' contributions

All authors contributed equally and significantly in writing this article. All authors read and approved the final manuscript.

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